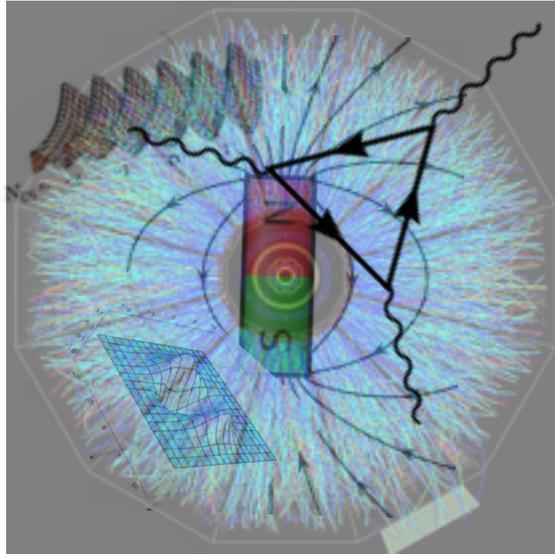


## Updates from CUJET3



**Jinfeng Liao**

Indiana University, Physics Dept. & CEEM

*Research Supported by NSF & DOE*



# Outline

- Jet (geometric) tomography, and how I came into this
- CUJET3: What It Is & What It Does
- HF Tests from CUJET3
- Summary & Discussions

## *References:*

*Shi, Xu, JL, Gyulassy, in preparation;*

*Xu, JL, Gyulassy, arXiv:1411.3673[CPL2015]; arXiv:1508.00552[JHEP2016].*

*Li, JL, Huang, PRD89,126006(2014).*

*X. Zhang, JL, PRC2013; PRC2014; PLB2012; arXiv:1311.5463.*

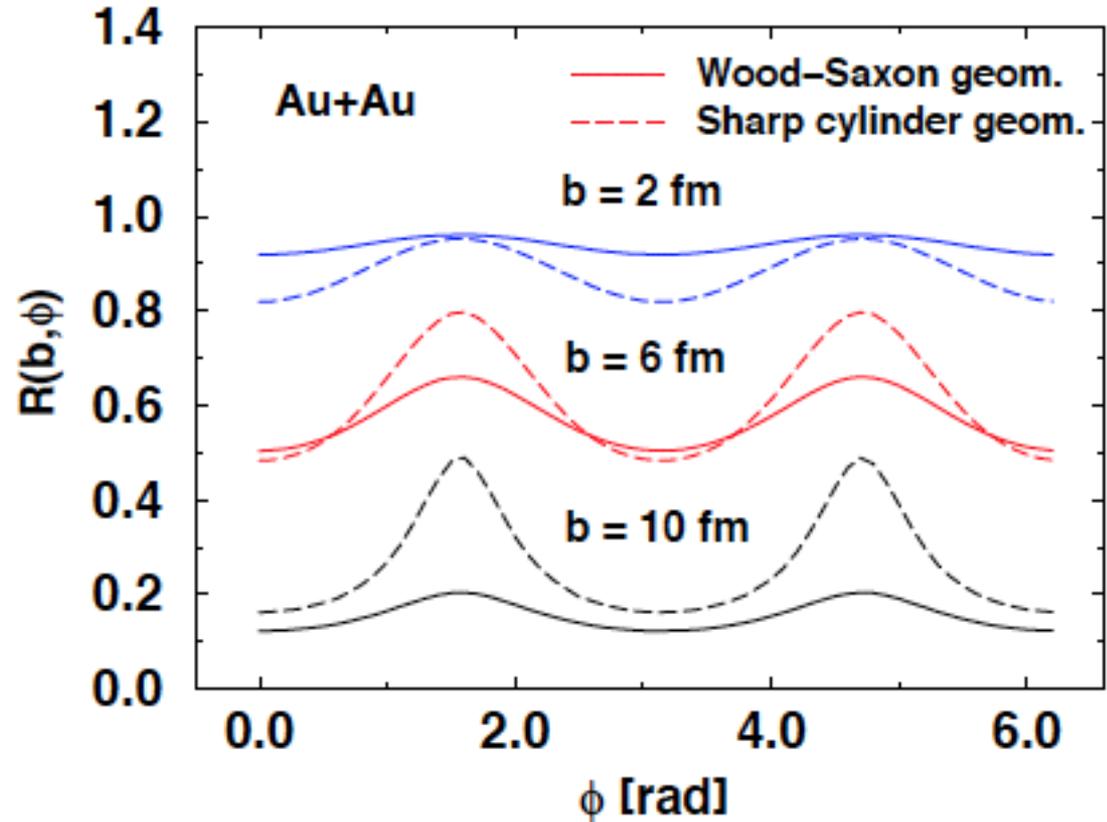
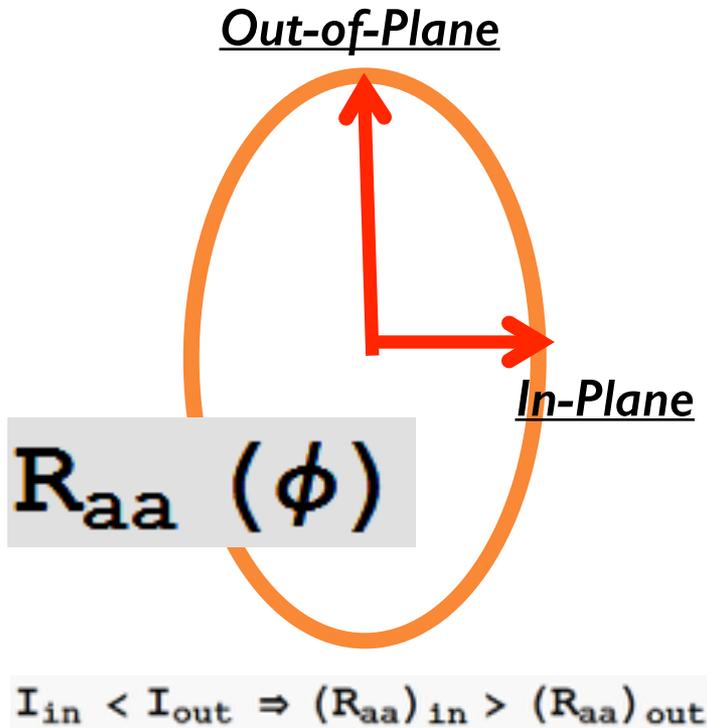
*JL, arXiv:1109.0271[PANIC11 proceedings].*

*JL, Shuryak, Phys.Rev.Lett. 102 (2009) 202302.*

# Geometric Anisotropy of Jet Quenching

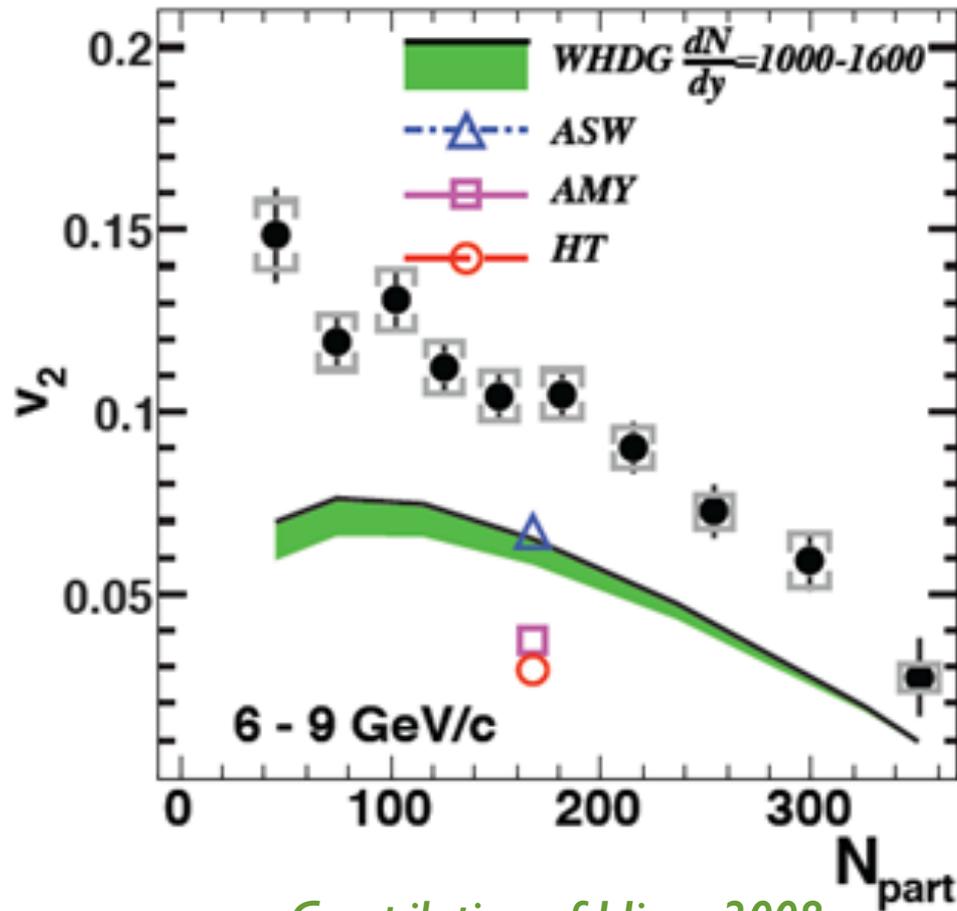
Geometric tomography (~2000)

[Gyulassy, Vitev, Wang, arXiv:nucl-th/0012092]

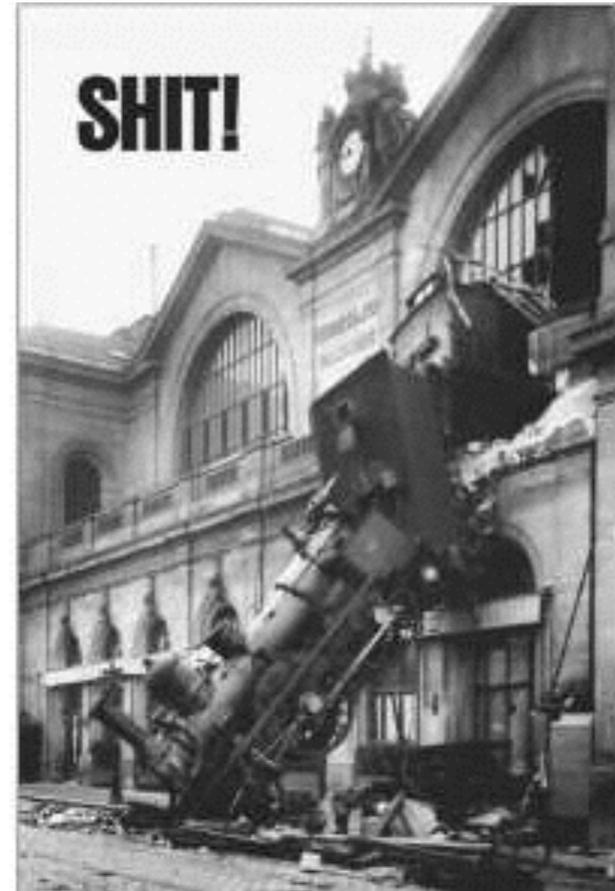


Positive  $v_2$  for high Pt hadrons — beautiful idea!  
It could be a “crowning” confirmation of jet energy loss model.

# Maybe Sometimes We Got Too Ambitious...



Compilation of J.Jia, ~2008



Till ~2008: clear and significant discrepancy between data / any model.  
That was the situation when I dived into this water...

# Magnetic Component of sQGP

## Magnetic Component of Strongly Coupled Quark-Gluon Plasma

A Dissertation Presented

by

Jinfeng Liao

to

The Graduate School

in Partial Fulfillment of the Requirements

for the Degree of

Doctor of Philosophy

in

Physics

Stony Brook University

August 2008

***There is the strongly coupled QGP,  
in the  $1\sim 2T_c$  regime.***

***What is the making of sQGP?***

***What is the connection to the  
vacuum confinement?***

***How the knowledge about sQGP  
helps us understand confinement?***

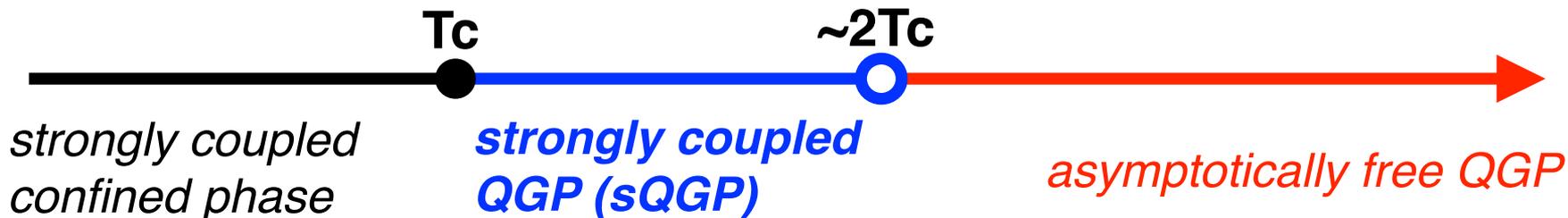
***“Magnetic Component”***

# sQGP: The Matter Just About to Be Confining

*The old dream*



*The new paradigm thanks to discoveries at RHIC and LHC ( $1\sim 3T_c$ ):*

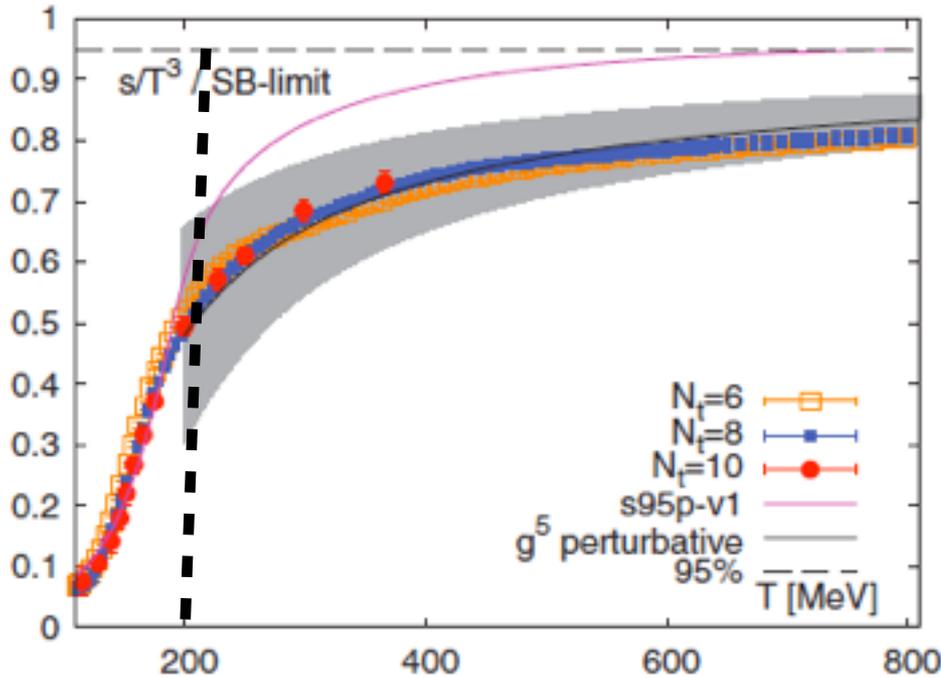


*A few of us, quickly came to realize:  
The matter just above confinement (in  $1\sim 2T_c$ ),  
is more like the confined world, rather than like  
the far-far-away place of asymptotic QGP!*

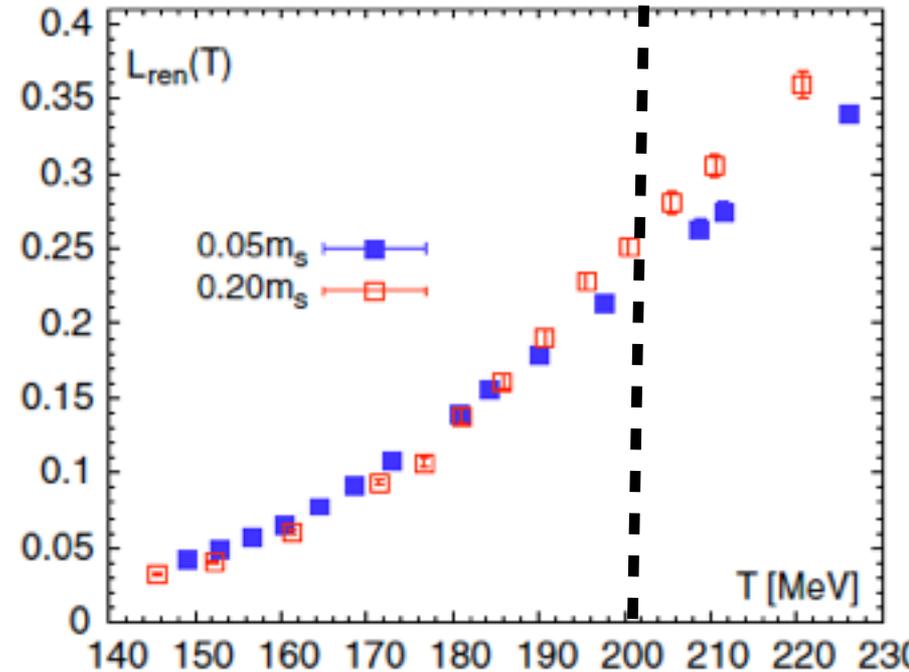
**This is to say, the confinement physics (whatever it is),  
must continue robustly into this region  
— we call it “postconfinement” regime!**

# The Special Near-Tc Matter: What are the DoF?

## Liberation of Thermal DoF



## Degree of color liberation



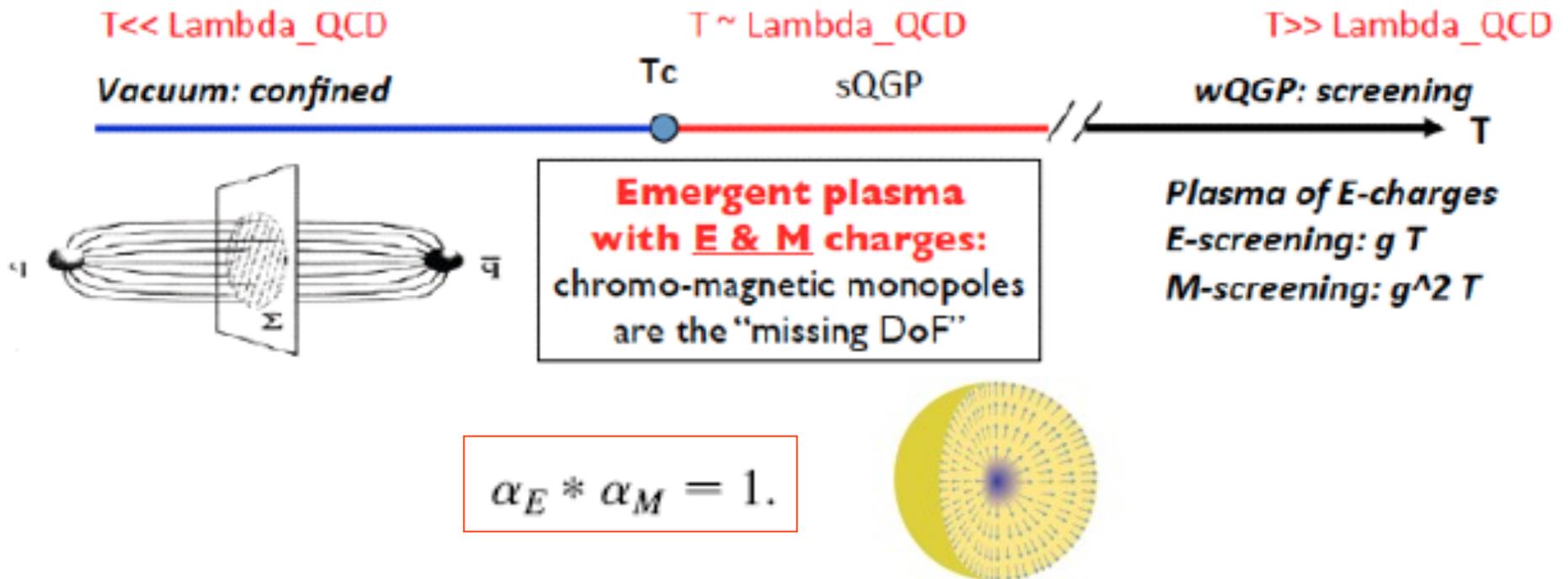
**Pisarski, ...: this is a semi-QGP!**

**Shuryak, Liao, ...: this is a monopole plasma!**

*The two pictures are in complement, from Electric or Magnetic language respectively, and reconciled into one coherent message:*

*Postconfinement Regime in  $1 \sim 2T_c$ ,  
with lively new manifestation of confinement physics*

# Magnetic Scenario of Near-Tc Plasma



**Condensate monopoles  $\rightarrow$  dense thermal monopoles  $1-2T_c$ : thermal monopoles hold together electric flux, yet with dissipation.**

PHYSICAL REVIEW C 75, 054907 (2007)

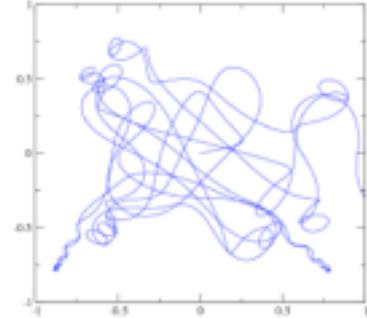
**Strongly coupled plasma with electric and magnetic charges**

Jinfeng Liao and Edward Shuryak

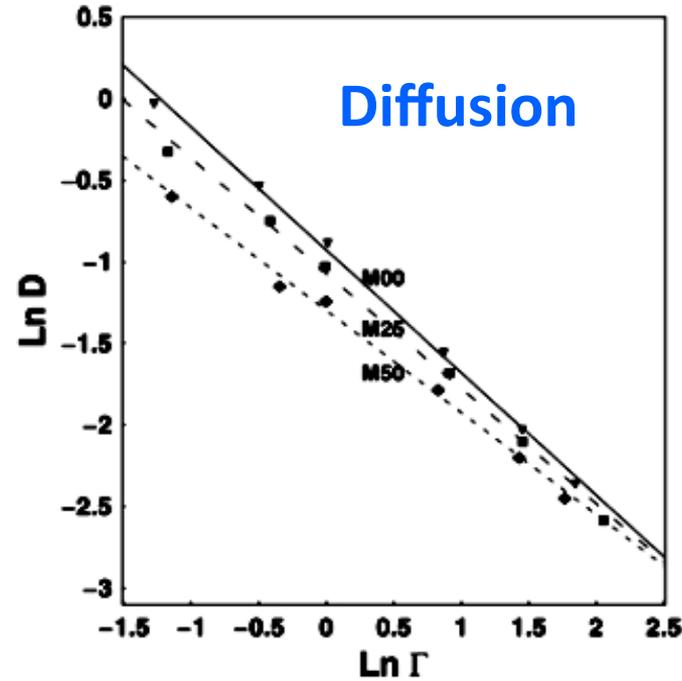
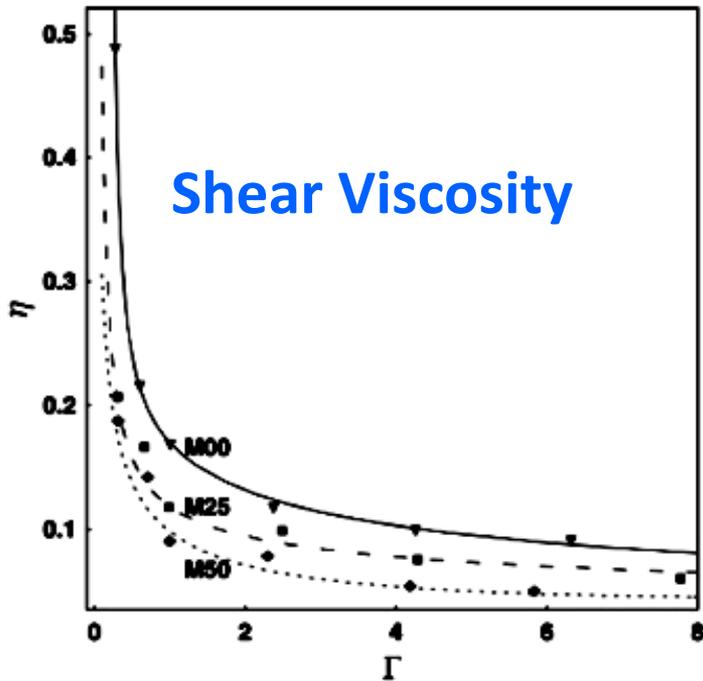
# WHAT A MAGNETIC COMPONENT DOES?

*We first studied the plasma of a completely new kind:  
Coulomb-Lorentz Plasma!*

Molecular Dynamics for 1000 particles with long range forces  
for varying E/M ratio:  
pure electric ; 25% magnetic charges ; 50% magnetic charges



*“Lorentz-Trapping Effect”*

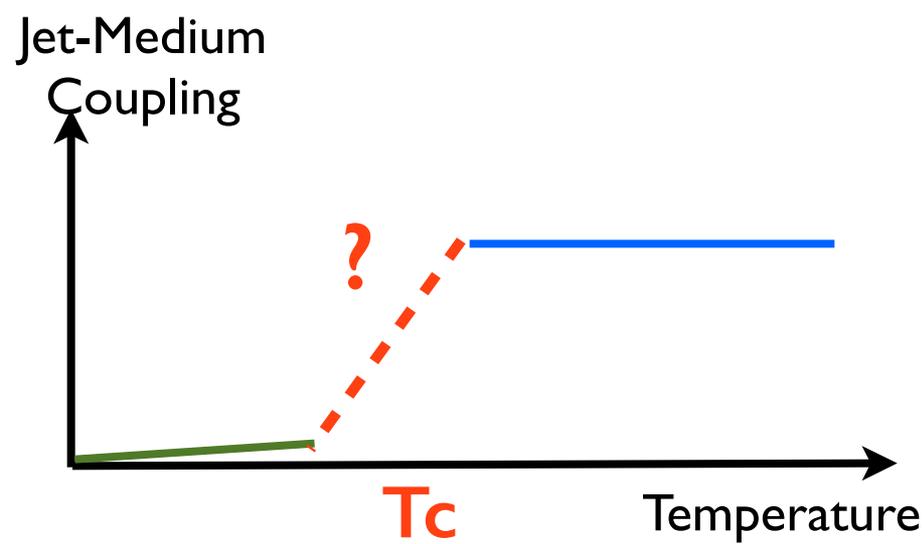


*A mixture of  
E&M charges  
can help  
explain the  
“perfect fluid”.*

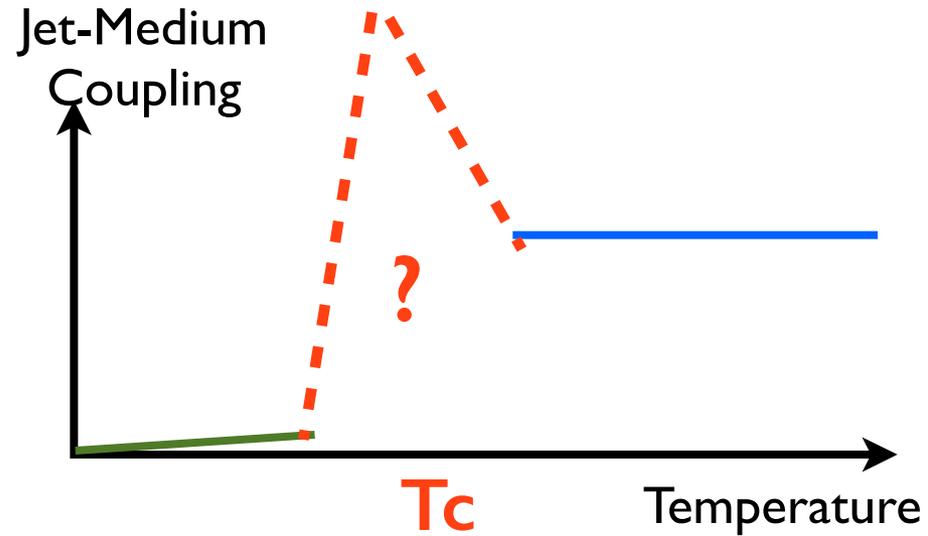
*[quantum E-M  
scattering:  
Ratti, Shuryak]*

**What about the jets? Can they sense the M charges?  
We spent a busy summer of 2008 on this problem...**

# From “Transparency” to Opaqueness



“Waterfall” scenario

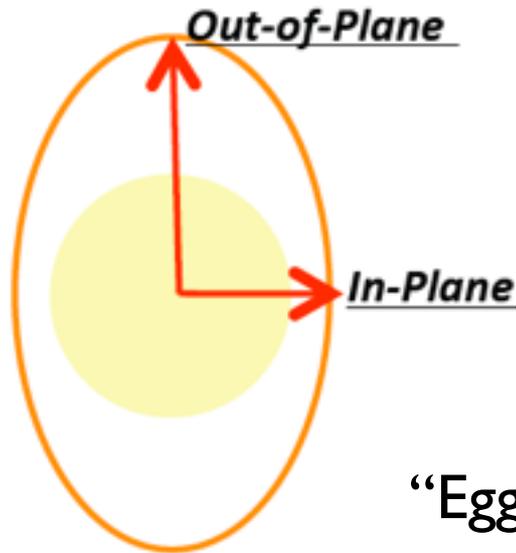


“Volcano” scenario



**The temperature dependence of jet-medium coupling has profound consequences!**

# Where Are Jets Quenched (More Strongly)?



Taken for granted in all previous models:  
“waterfall” scenario.

We realized the puzzle may concern  
more radical questions:

Where are jets quenched (more strongly)?

Geometry is a sensitive feature:  
“Egg yolk” has one geometry, “Egg white” has another.

## Where are jets quenched in heavy-ion collisions?

Jinfeng Liao<sup>1,2\*</sup> and Edward Shuryak<sup>1†</sup>

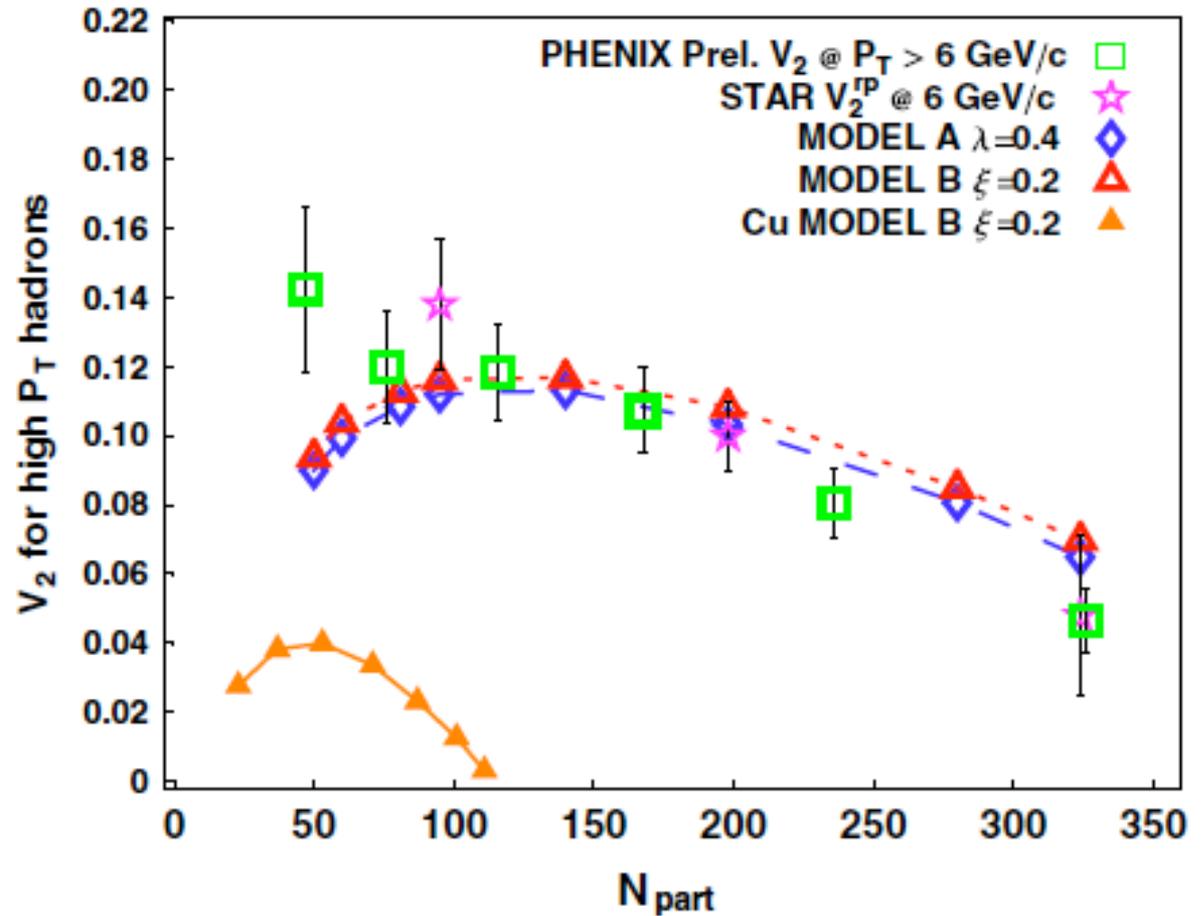
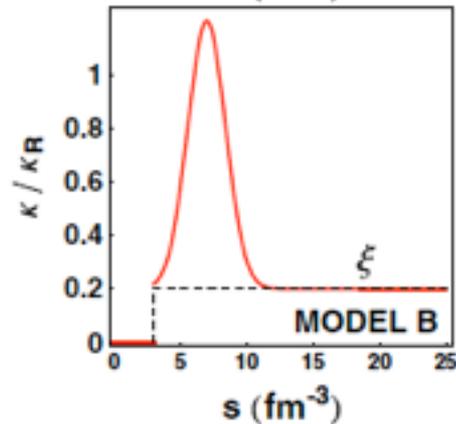
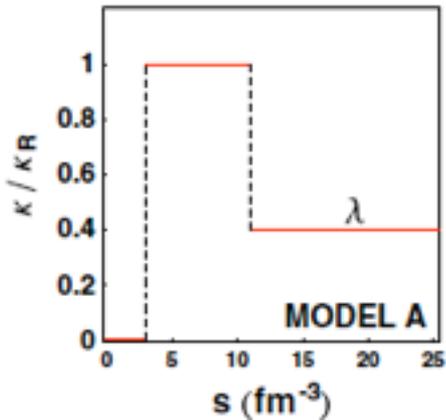
<sup>1</sup>*Department of Physics and Astronomy, State University of New York, Stony Brook, NY 11794*

<sup>2</sup>*Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720*

(Dated: October 22, 2008)

We study dependence of jet quenching on matter density, using “tomography” of the fireball provided by RHIC data on jet azimuthal asymmetry parameter  $v_2(b)$  for large  $p_t$  hadrons. Slicing the fireball into shells with constant (entropy) density, we derive a new geometrical limit for it which indeed is above the data  $v_2(b) < v_2^{max}(b)$ . Interestingly, the limit is reached only if quenching is dominated by a shell with the entropy density  $3 fm^{-3} < s < 11 fm^{-3}$ , exactly the near- $T_c$  region. We conclude that the data can be explained if quenching is few times stronger in the near- $T_c$  region than in the QGP at  $T > T_c$ . We also argue that recent views picturing the near- $T_c$  region as a magnetic plasma of monopoles can naturally explain that.

# Near-Tc Enhancement (NTcE)

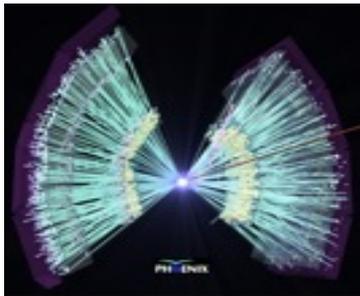
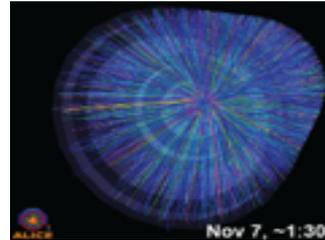
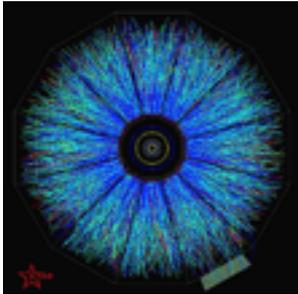


In the paper we concluded:

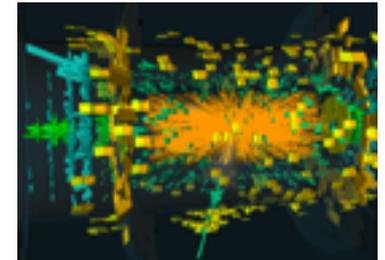
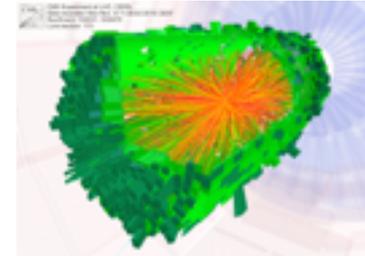
“In relativistic heavy ion collisions the jets are quenched about **2--5 times stronger** in the near-Tc region than the higher-T QGP phase.”

# Hot QCD Matter from RHIC to LHC

## RHIC Events



## LHC Events



**From RHIC to LHC:  
Capability of shifting  
QGP Temperature!**

## The Surprising Transparency of the sQGP at LHC

W. A. Horowitz\*

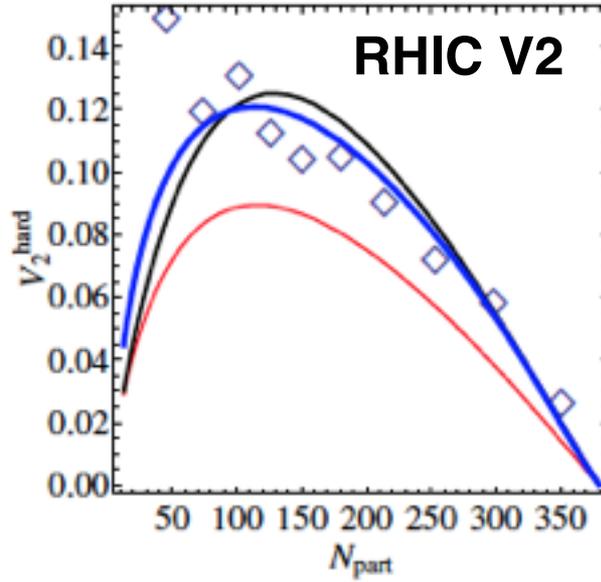
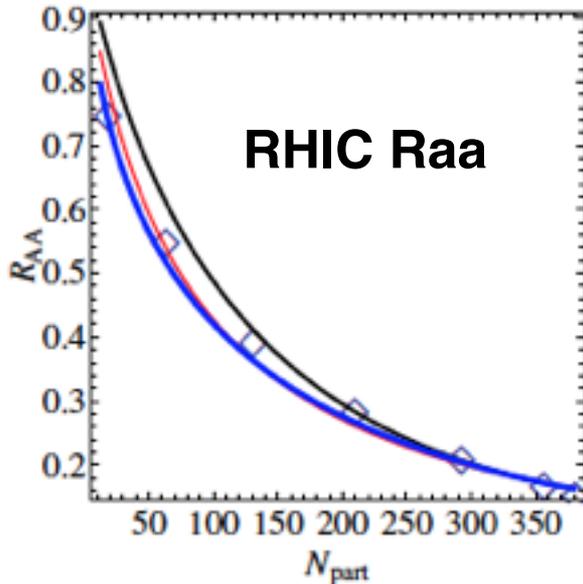
*Department of Physics, University of Cape Town,  
Private Bag X3, Rondebosch 7701, South Africa*

Miklos Gyulassy

*Department of Physics, Columbia University,  
538 West 120<sup>th</sup> Street, New York, NY 10027, USA*

(Dated: April 27, 2011)

# Jet Tomo vs Mono vs Holo -graphy



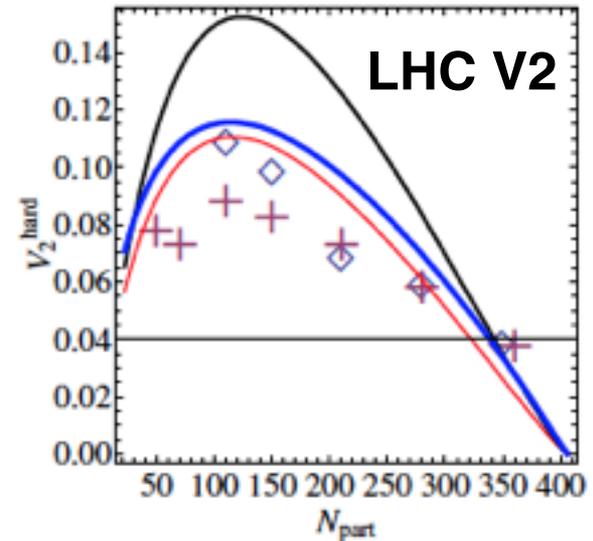
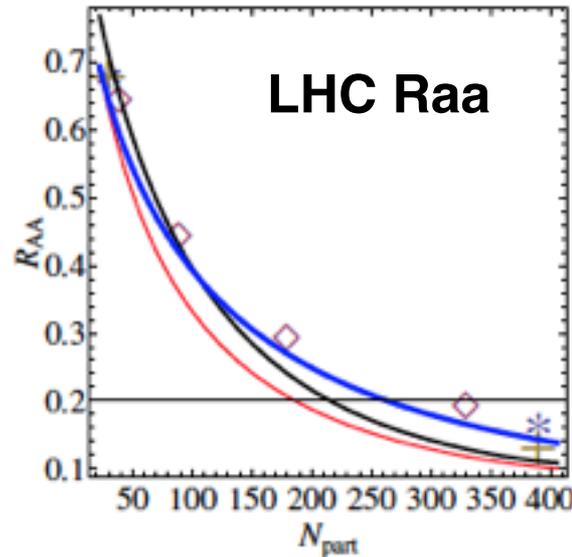
*My calculation results were reported at PANIC2011, July @MIT*

**RED:  $L^2$  model**

**BLUE:  $L^2$  + Near-Tc**

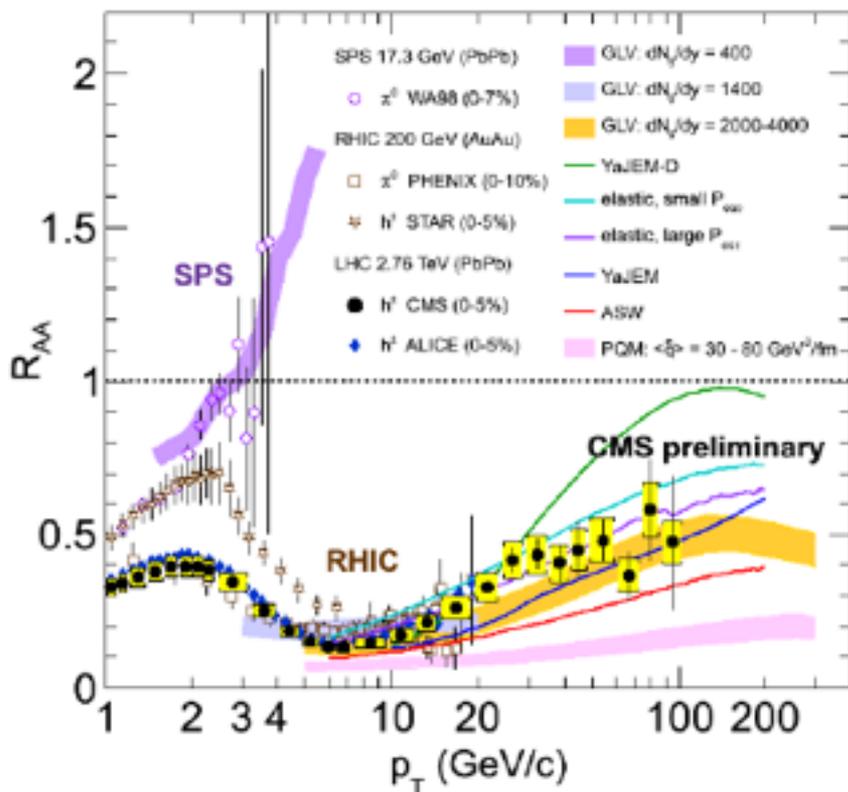
**BLACK:  $L^3$  model**

*Liao, "Nonperturbative Jet Quenching from Geometric Data" arXiv:1109.0271[nucl-th]*

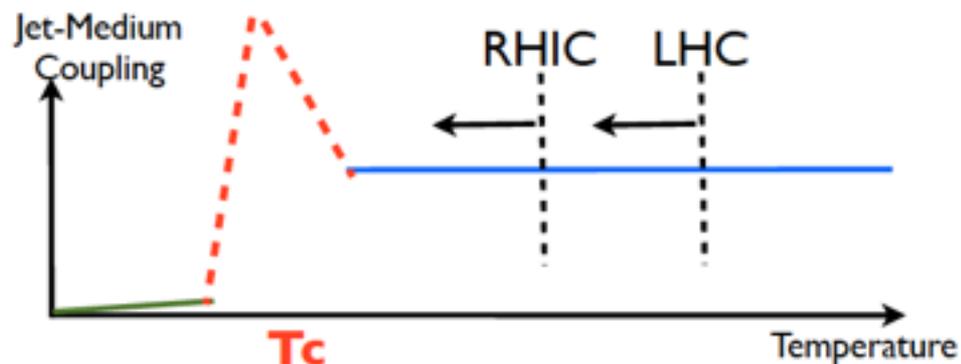


# The Opaqueness Does Shift!

Beautiful jet quenching measurements  
from all collaborations



Already a clear hint of LESS OPACITY:  
similar  $R_{AA}$ , despite twice the density!  
— “surprising transparency” (Horowitz  
& Gyulassy, QM I I)  
— naturally expected if the “volcano  
scenario” is indeed true (Liao PANIC I I)



# Quantifying the Reduction of Opaqueness

Examining a reduced jet-medium coupling in Pb+Pb collisions at the Large Hadron Collider

Barbara Betz<sup>a</sup> and Miklos Gyulassy<sup>b</sup>

<sup>a</sup>*Institute for Theoretical Physics, Johann Wolfgang Goethe-University, 60438 Frankfurt am Main, Germany*

<sup>b</sup>*Department of Physics, Columbia University, New York, 10027, USA*

Effective Coupling $\kappa$ assuming $\tau_0 = 0.01$ fm/c			
$\sqrt{s}$	Glauber z=1	dcgc1.2 z=1	Glauber z=2
0.20	0.60	0.58	0.44
2.76	0.45	0.43	0.26
LHC/RHIC	0.75	0.74	0.59

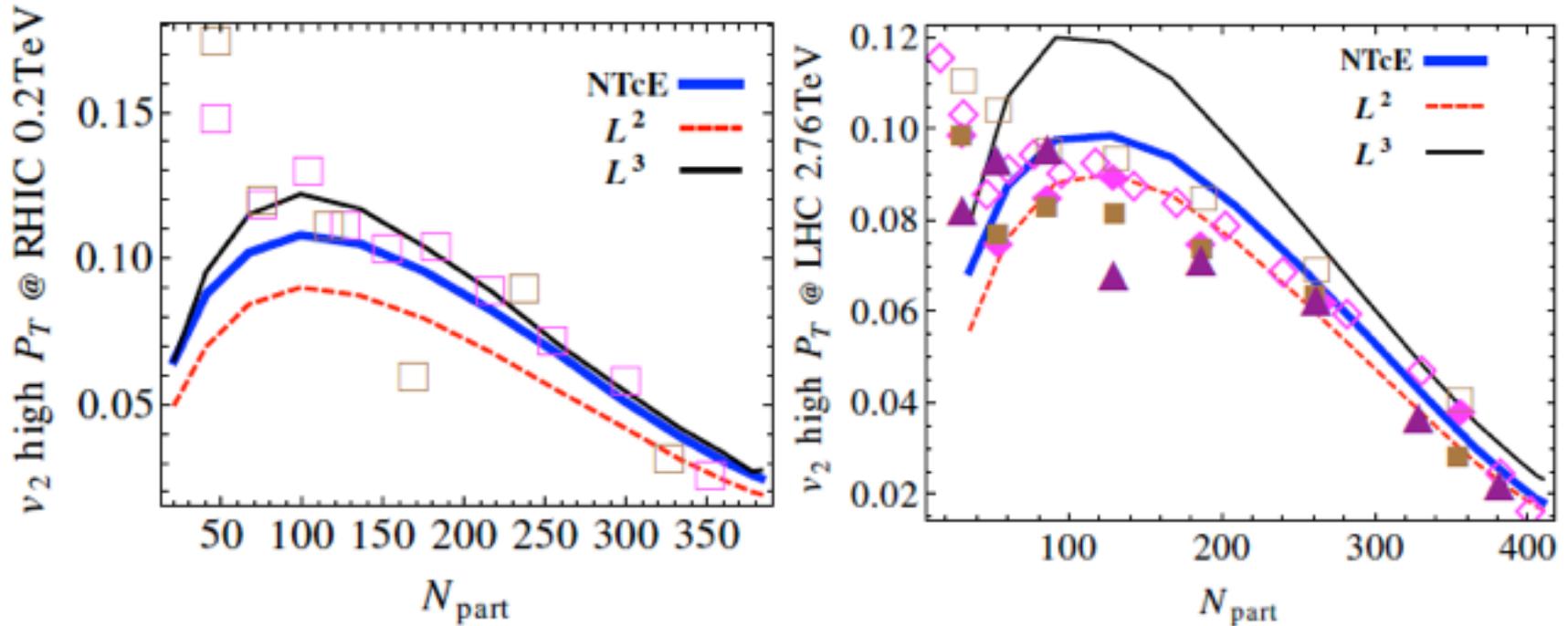
Analysis by me and student came to the same conclusion.

*Zhang & JL, arXiv: 1208.6361, 1210.1245.*

$$\langle \kappa \rangle_{\text{RHIC}} : \langle \kappa \rangle_{\text{LHC}} \approx 1 : 0.72$$

Let me emphasize: this reduction is naturally born out from near- $T_c$  enhancement !

# V2 from RHIC to LHC



RED:  $L^2$  model+waterfall

BLUE:  $L^2$ +volcano

BLACK:  $L^3$ +waterfall

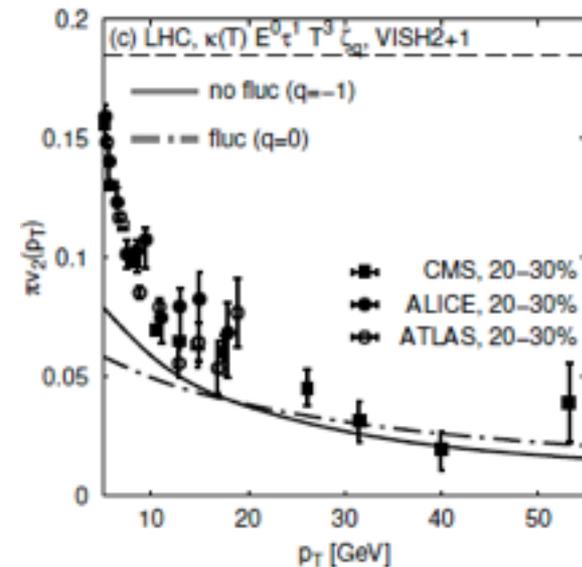
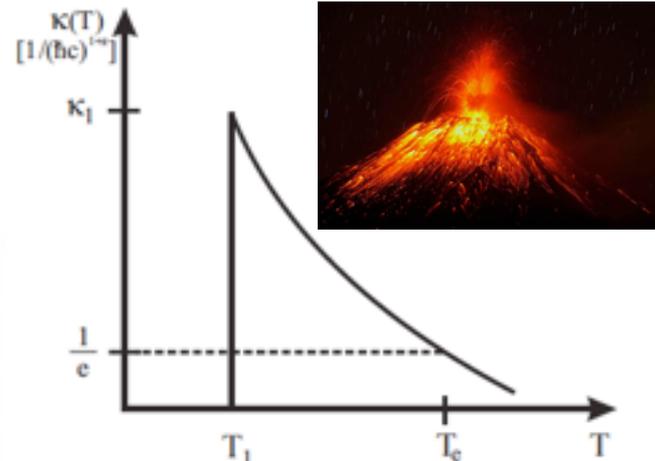
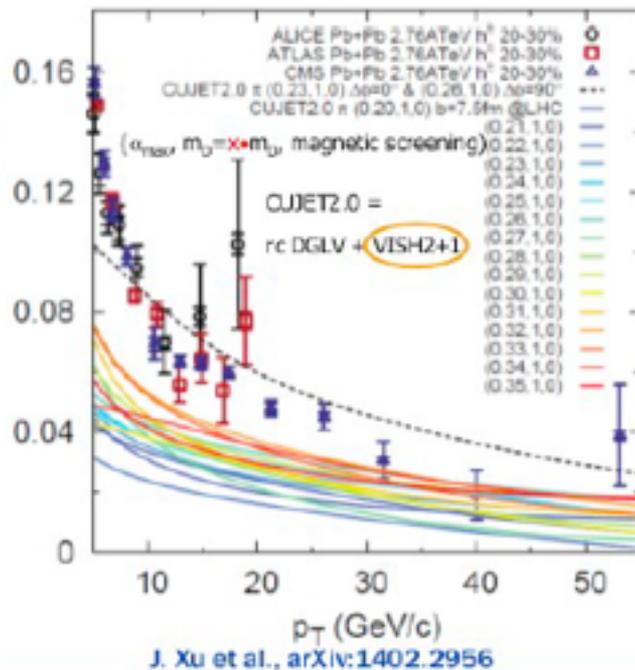
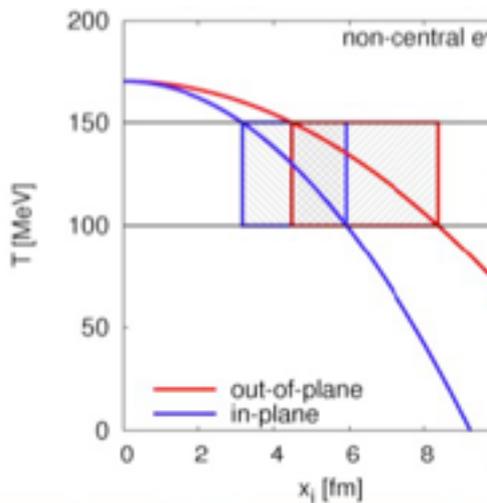
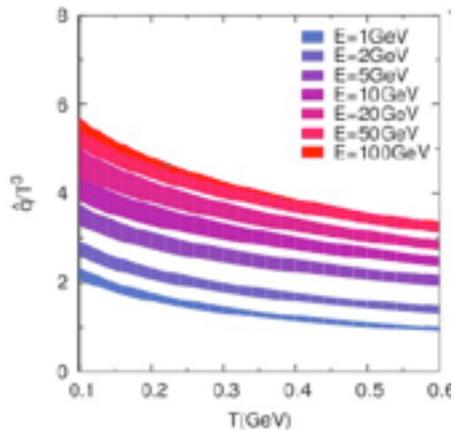
\* We do see big difference between waterfall/volcano at RHIC, and this difference becomes much smaller at LHC

\* RHIC + LHC data are in favor of the  $L^2$  + Volcano scenario

Zhang & JL, arXiv: 1208.6361

# Detailed Analysis from Earlier CUJET

Xu, Buzzatti, Gyulassy, 1402.2956;  
Betz, Gyulassy, 1404.6378

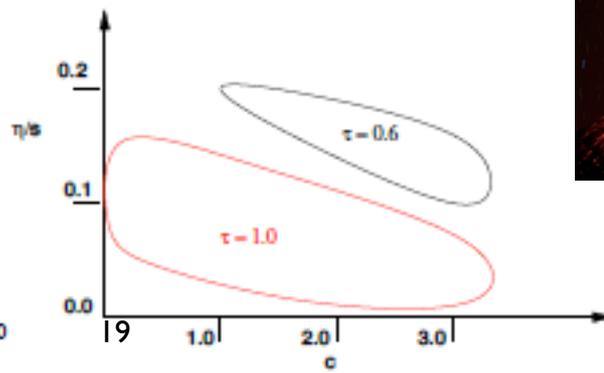
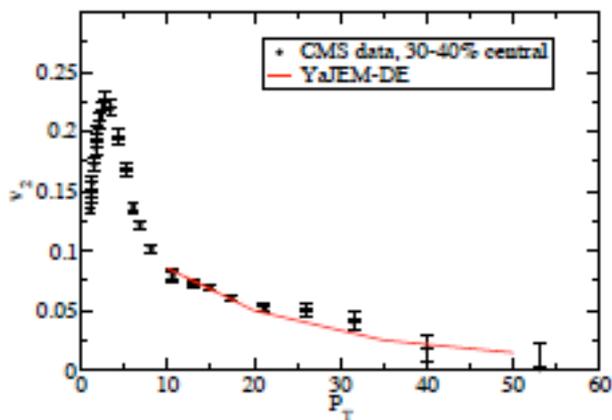
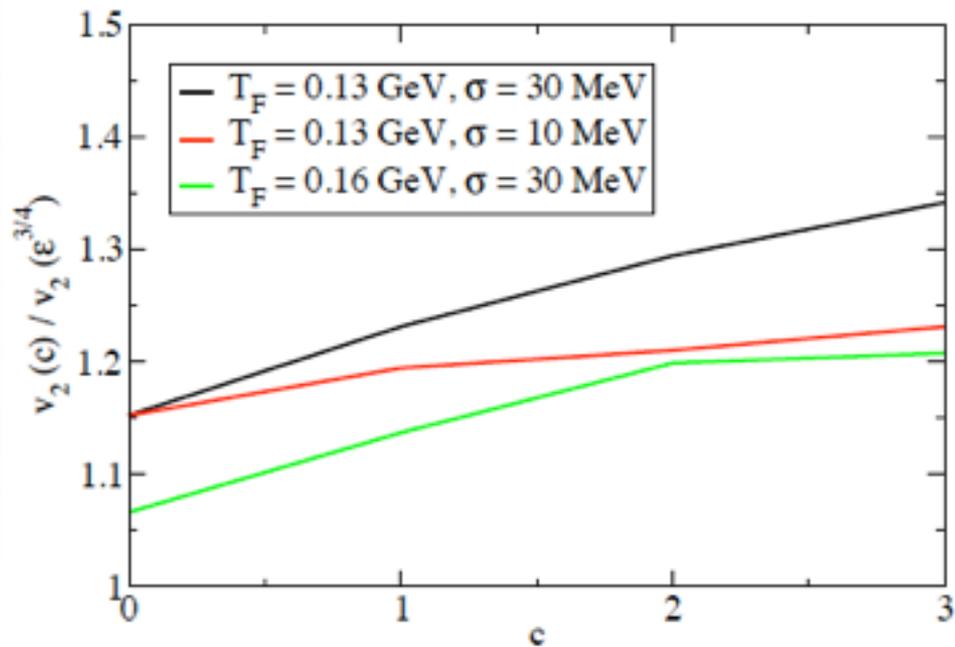
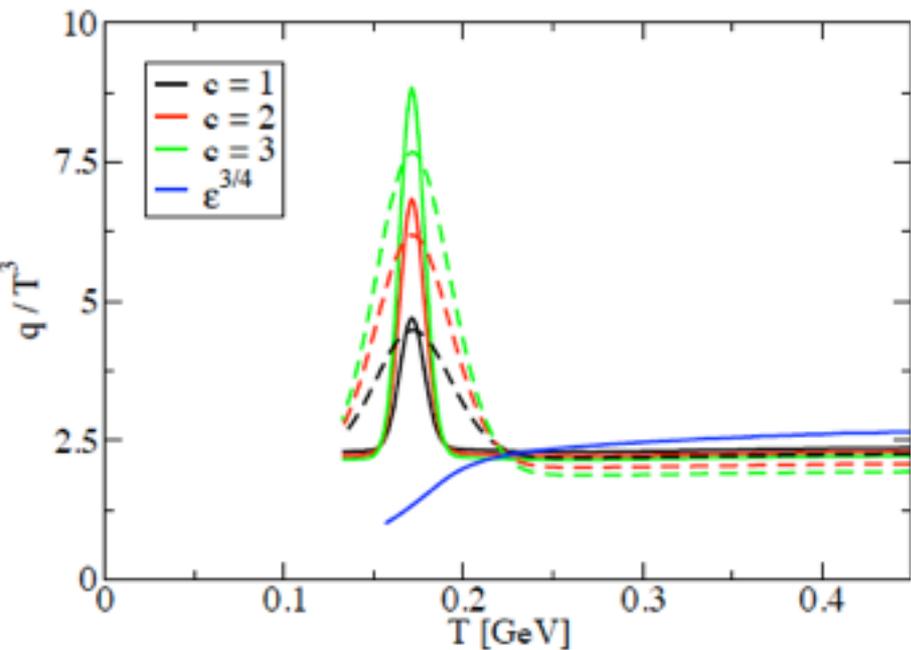


effectively:  
 $\kappa_{out} > \kappa_{in}$

# Results from Renk's Simulations

Renk, 1402.5798 & QM14

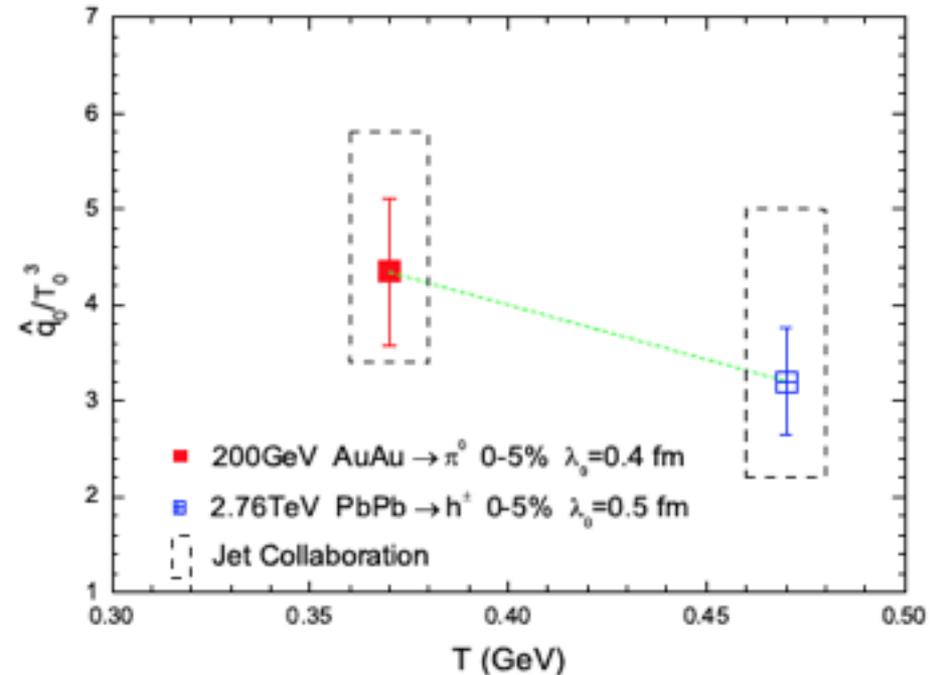
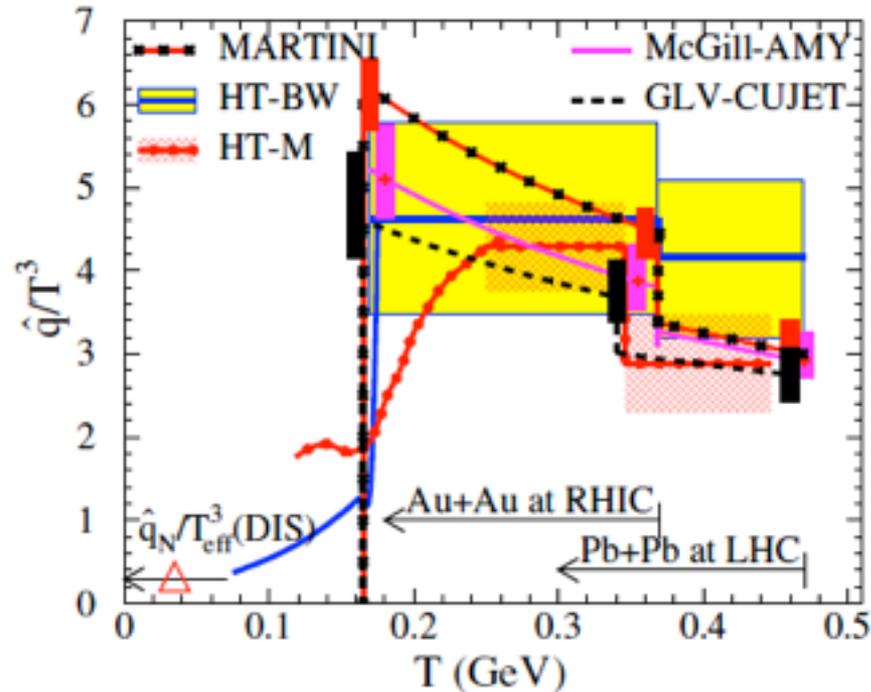
model	ASW	YDE 3d	YDE 2d
NTC/ $\epsilon^{3/4}$	1.17	1.22	1.20



# Extractions at RHIC vs LHC

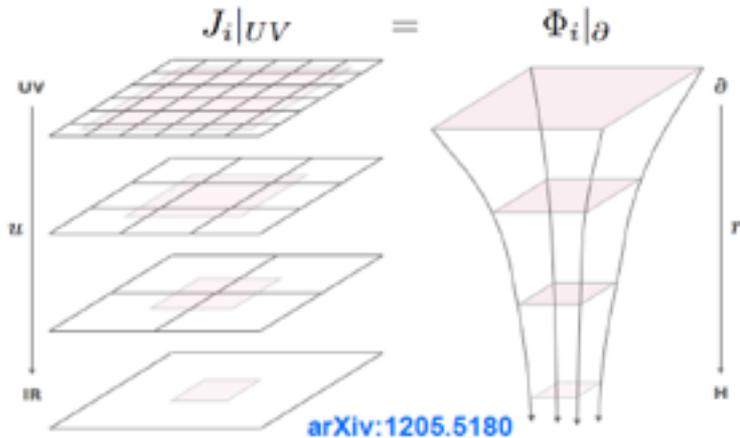
JET Collaboration,  
arXiv: 1312.5003

Liu, Zhang, Zhang, Wang,  
arXiv: 1506.02840



In the paper PRL(2009) we (Liao&Shuryak) concluded:  
“In relativistic heavy ion collisions the jets are quenched about  
**2--5 times stronger** in the near- $T_c$  region  
than the higher- $T$  QGP phase.”

# Message from One More Dimension

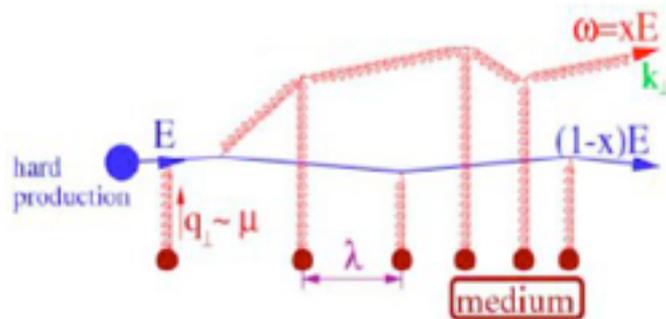


Deforming the conformal-AdS to introduce non-conformal dynamics: using graviton-dilaton system in the bulk

$$S_G = \frac{1}{16\pi G_5} \int d^5x \sqrt{g_s} e^{-2\Phi} (R_s + 4\partial_M \Phi \partial^M \Phi - V_G^s(\Phi))$$

$$\Phi(z) = \mu_G^2 z^2 \tanh(\mu_G^4 z^2 / \mu_G^2)$$

$$ds_S^2 = e^{2A_s} \left( -f(z) dt^2 + \frac{dz^2}{f(z)} + dx^i dx^i \right)$$

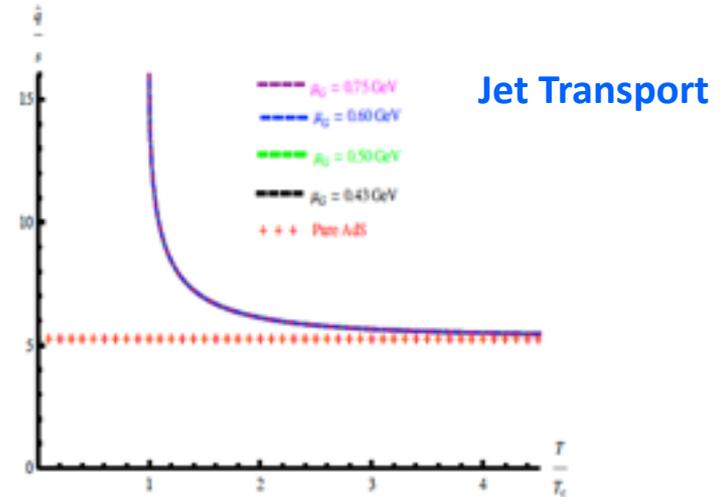
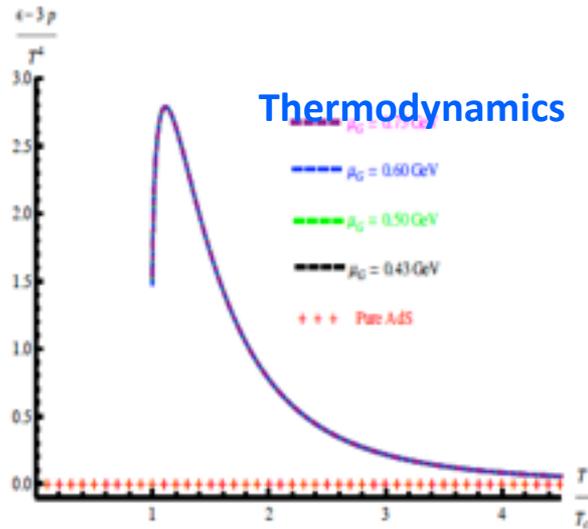


We use the Liu-Rajagopal-Wiedemann scheme to compute  $\hat{q}$

$$\hat{q} = \frac{\sqrt{2}\sqrt{\lambda}}{\pi \int_0^{z_h} dz \sqrt{g_{zz} / (g_{22}^2 g_{--})}},$$

D. Li, JL, M. Huang, arXiv:1401.2035

# Results from Non-Conformal Holo-QCD

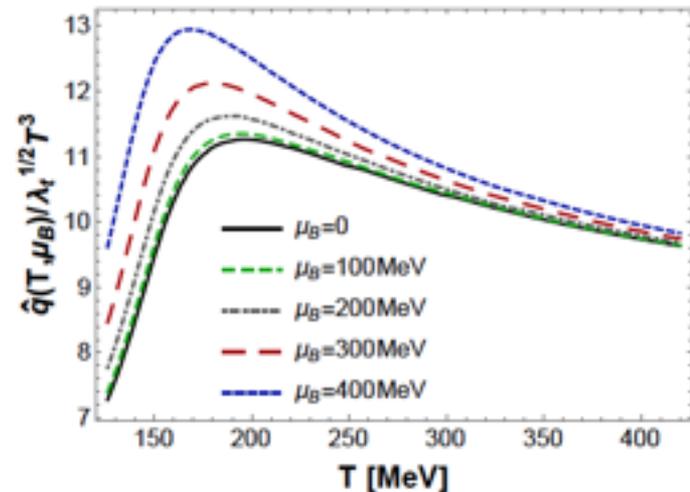


*D. Li, JL, M. Huang, arXiv:1401.2035*

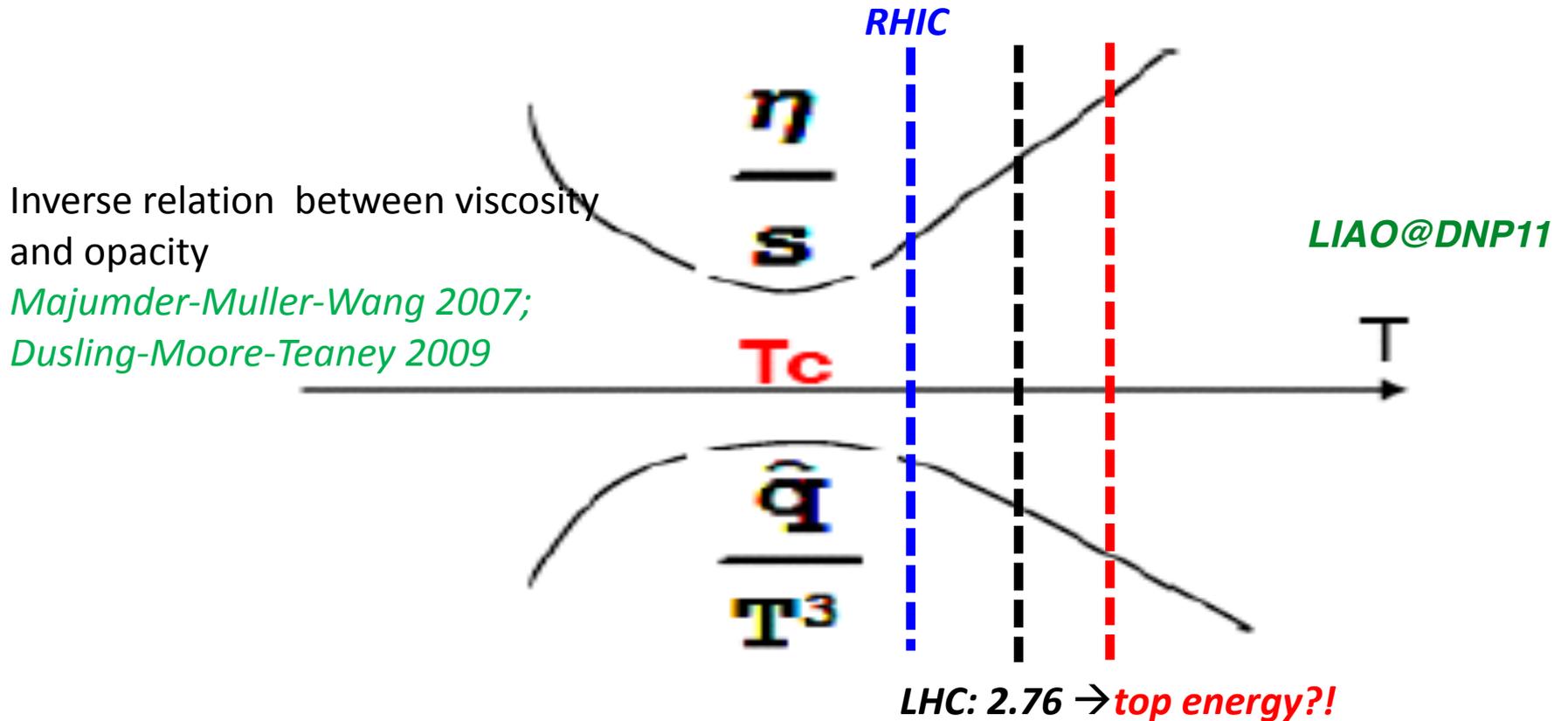
**Same non-conformal, non-monotonic, non-perturbative dynamics  
 ---> shows up in trace anomaly and in jet transport parameter**

Rougemont, Ficnar, Finazzo, Noronha,  
 arXiv:1507.06556

Quite different holo setup, but showing  
 the same robust connection as above!



# NEAR-Tc MATTER IS SPECIAL



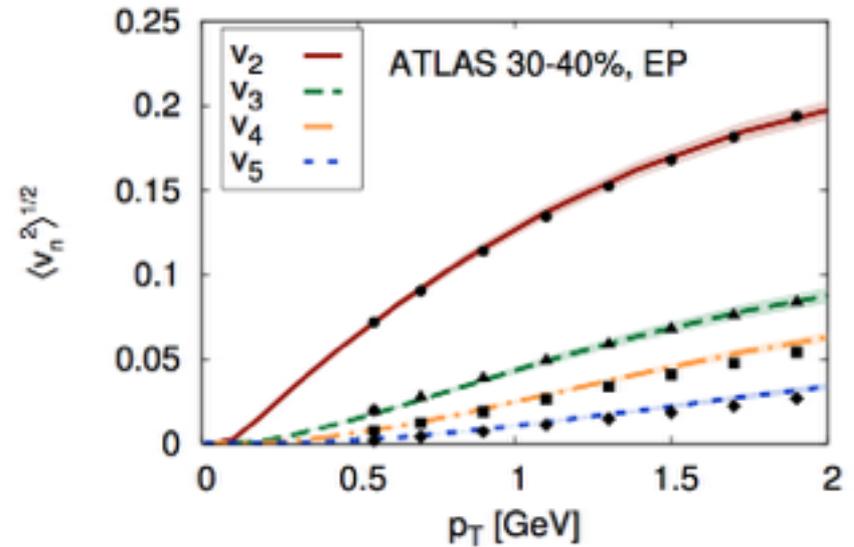
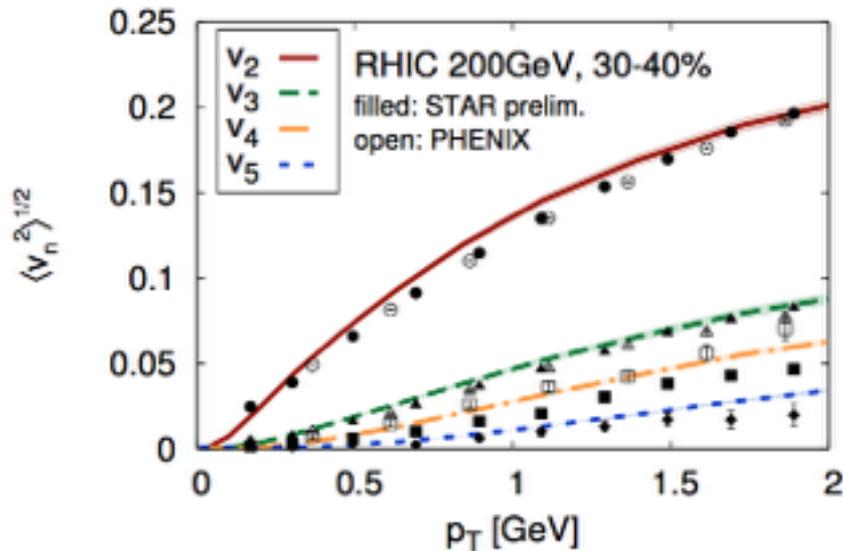
Will we see a systematic deviation from RHIC to LHC? Yes!

**The “see-saw”-QGP expects such a picture to occur in a narrow regime  $1-3T_c$ .**

**A kind of “critical opalescence”!**

**Reminiscence of a phase transition underlying the crossover**

# The QGP Liquidity is Shifting Too!



$$\eta/s \approx 0.12 \longrightarrow \eta/s \approx 0.2$$

Works of multiple groups (BNL-McGill, Frankfurt, Scalay, OSU) consistently suggest a visible increase,  $\sim 40\%$ , of average  $\eta/s$  from RHIC to LHC.

*To be in context: the temperature is increased only by  $\sim 30\%$  from RHIC to LHC.*

*Such rapid change is an indication of near- $T_c$  phenomenon.*

# Toward Microscopic Making of sQGP!

*There are a number of outstanding challenges in understanding how the QGP does what it does:*

- \* We know that there are nonperturbative dynamics and emergent degrees of freedom in sQGP — how to implement such physics?*
- \* Experimental & lattice data validation/constraints?*
- \* Perfect fluidity v.s. Jet quenching — how to reconcile the two key properties of the sQGP?*



*Xu, JL, Gyulassy,  
arXiv:1411.3673;  
arXiv:1508.00552*

# CUJET3: Semi-Quark-Gluon Monopole Plasma

CHIN PHYS LETT Vol 32, No. 9 (2015) 092501

Express Letter

## Consistency of Perfect Fluidity and Jet Quenching in Semi-Quark-Gluon Monopole Plasmas \*

Jiechen Xu(徐杰晨)<sup>1</sup>, Jinfeng Liao(廖劲峰)<sup>2,3\*\*</sup>, Miklos Gyulassy<sup>1\*\*</sup>

<sup>1</sup>Department of Physics, Columbia University, New York 10027, USA

<sup>2</sup>Physics Department and CEFM, Indiana University, Bloomington 47408, USA

<sup>3</sup>RIKEN BNL Research Center, Bldg. 510A, Brookhaven National Laboratory, New York 11973, USA

(Received 31 July 2015)

We utilize a new framework, CUJET3.0, to deduce the energy and temperature dependence of the jet transport parameter,  $\hat{q}(E > 10 \text{ GeV}, T)$ , from a combined analysis of available data on nuclear modification factor and azimuthal asymmetries from high energy nuclear collisions at RHIC/BNL and LHC/CERN. Extending a previous perturbative-QCD based jet energy loss model (known as CUJET2.0) with (2+1)D viscous hydrodynamic bulk evolution, this new framework includes three novel features of nonperturbative physics origin: (i) the Polyakov loop suppression of color electric scattering (aka 'semi-QGP' of Pisarski et al.), (ii) the enhancement of jet scattering due to emergent magnetic monopoles near  $T_c$  (aka 'magnetic scenario' of Liao and Shuryak), and (iii) thermodynamic properties constrained by lattice QCD data. CUJET3.0 reduces to v2.0 at high temperatures  $T > 400 \text{ MeV}$ , while greatly enhances  $\hat{q}$  near the QCD deconfinement transition temperature range. This enhancement accounts well for the observed elliptic harmonics of jets with  $p_T > 10 \text{ GeV}$ . Extrapolating our data-constrained  $\hat{q}$  down to thermal energy scales,  $E \sim 2 \text{ GeV}$ , we find for the first time a remarkable consistency between high energy jet quenching and bulk perfect fluidity with  $\eta/s \sim T^2/\hat{q} \sim 0.1$  near  $T_c$ .

PACS: 25.75.-q, 12.38.Mh, 24.85.+p, 13.87.-a DOI: 10.1088/0256-307X/32/9/092501

## Bridging soft-hard transport properties of quark-gluon plasmas with CUJET3.0

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538 West 120th Street, New York, NY 10027, U.S.A.

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# A Sophisticated Simulation Framework

DGLV-CUJET framework for describing multi-parton scattering:

$$\begin{aligned}
 x_E \frac{dN_g^{n=1}}{dx_E} &= \frac{18C_R}{\pi^2} \frac{4 + N_f}{16 + 9N_f} \int d\tau n(\mathbf{z}) \Gamma(\mathbf{z}) \int d^2k \\
 &\times \alpha_s \left( \frac{\mathbf{k}^2}{x_+(1-x_+)} \right) \int d^2q \frac{\alpha_s^2(\mathbf{q}^2)}{\mu^2(\mathbf{z})} \frac{f_E^2 \mu^2(\mathbf{z})}{\mathbf{q}^2(\mathbf{q}^2 + f_E^2 \mu^2(\mathbf{z}))} \\
 &\times \frac{-2(\mathbf{k} - \mathbf{q})}{(\mathbf{k} - \mathbf{q})^2 + \chi^2(\mathbf{z})} \left[ \frac{\mathbf{k}}{\mathbf{k}^2 + \chi^2(\mathbf{z})} - \frac{(\mathbf{k} - \mathbf{q})}{(\mathbf{k} - \mathbf{q})^2 + \chi^2(\mathbf{z})} \right] \\
 &\times \left[ 1 - \cos \left( \frac{(\mathbf{k} - \mathbf{q})^2 + \chi^2(\mathbf{z})}{2x_+ E} \tau \right) \right] \left( \frac{x_E}{x_+} \right) \left| \frac{dx_+}{dx_E} \right| \cdot (
 \end{aligned}$$

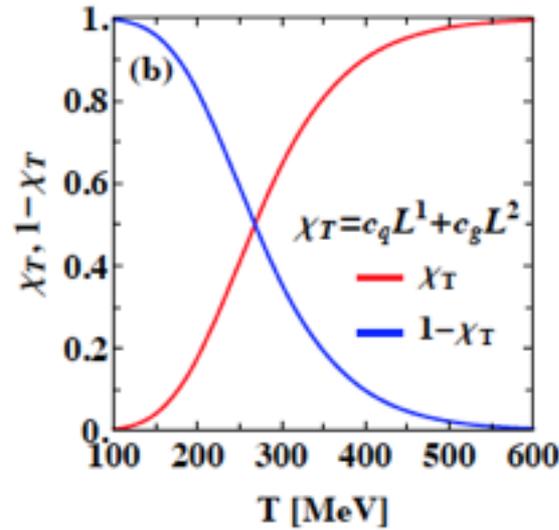
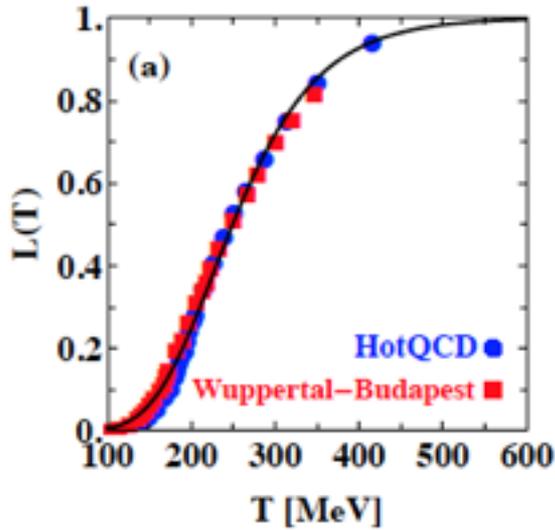
Original DGLV formalism has only quark/gluon scattering centers

We now include both color-electric and color-magnetic scattering centers.

$$x \frac{dN}{dx} \propto \dots \int_{q^2} \left[ \frac{n \alpha_s^2(q^2) f_E^2}{q^2(q^2 + f_E^2 \mu^2)} \right] \dots \longrightarrow \left[ \frac{n_e (\alpha_s(q^2) \alpha_s(q^2)) f_E^2}{q^2(q^2 + f_E^2 \mu^2)} + \frac{n_m (\alpha^e(q^2) \alpha^m(q^2)) f_M^2}{q^2(q^2 + f_M^2 \mu^2)} \right]$$

Our goal is to implement the nonperturbative NEAR-Tc Physics  
 → CUJET3.0

# The Making of sQGP

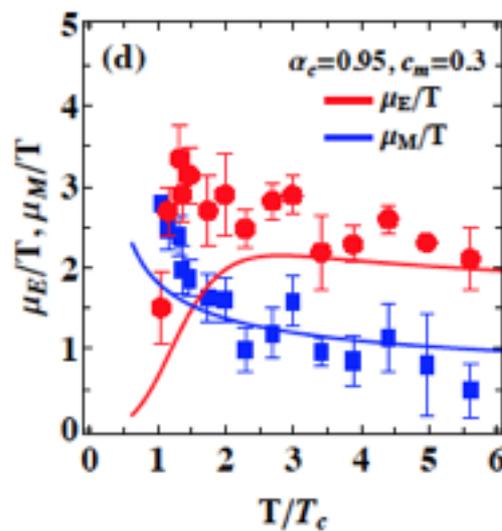
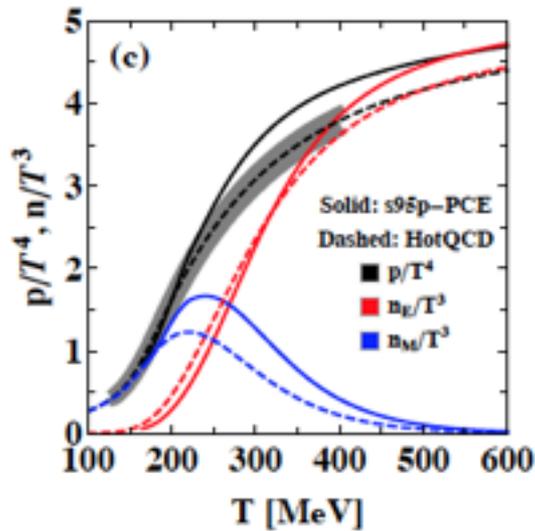


\* *Electric density:  
L-loop suppression*

$$\chi_T = c_q L + c_g L^2$$

\* *Magnetic density:  
constrained by total pressure*

$$(1 - \chi_T)$$



\* *Running coupling:  
L-loop suppression*

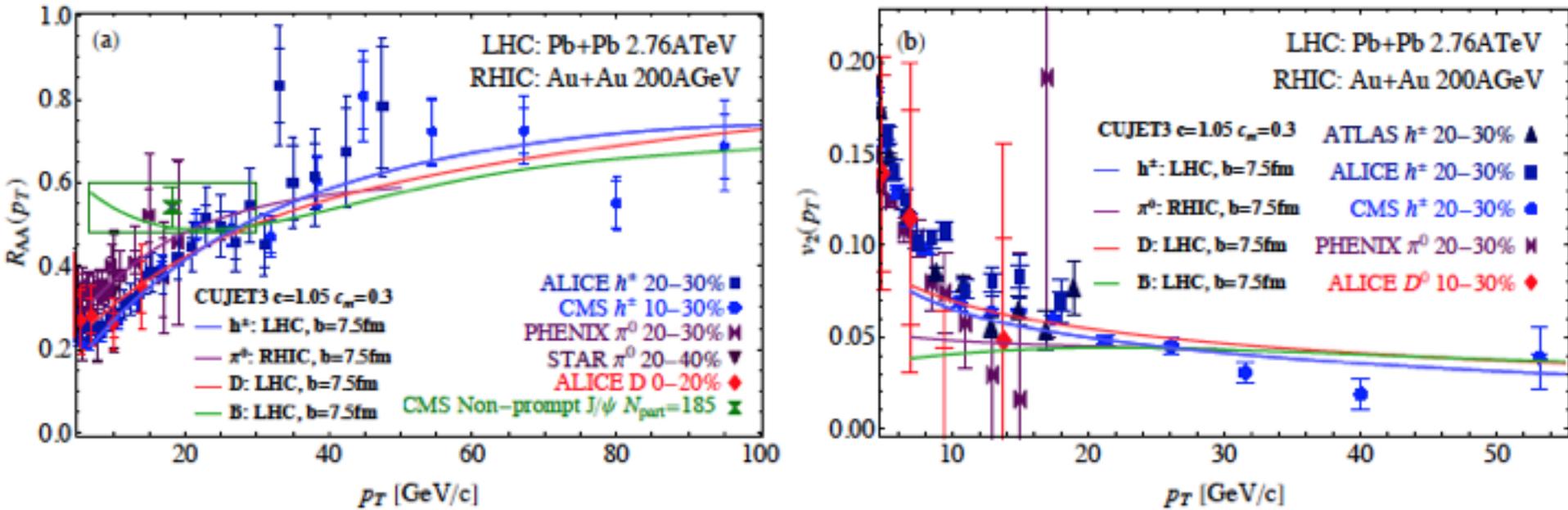
$$\alpha_s(Q^2) = \alpha_c / \left[ 1 + \frac{9\alpha_c}{4\pi} \text{Log}\left(\frac{Q^2}{T_c^2}\right) \right]$$

\* *Screening:*

$$f_E(T) = \sqrt{\chi_T} \quad , \quad f_M(T) = c_m g$$

The model implementations of electric and magnetic components are well constrained by available lattice data.

# CUJET3.0 Explains (RHIC+LHC)\*(Raa+V2)!

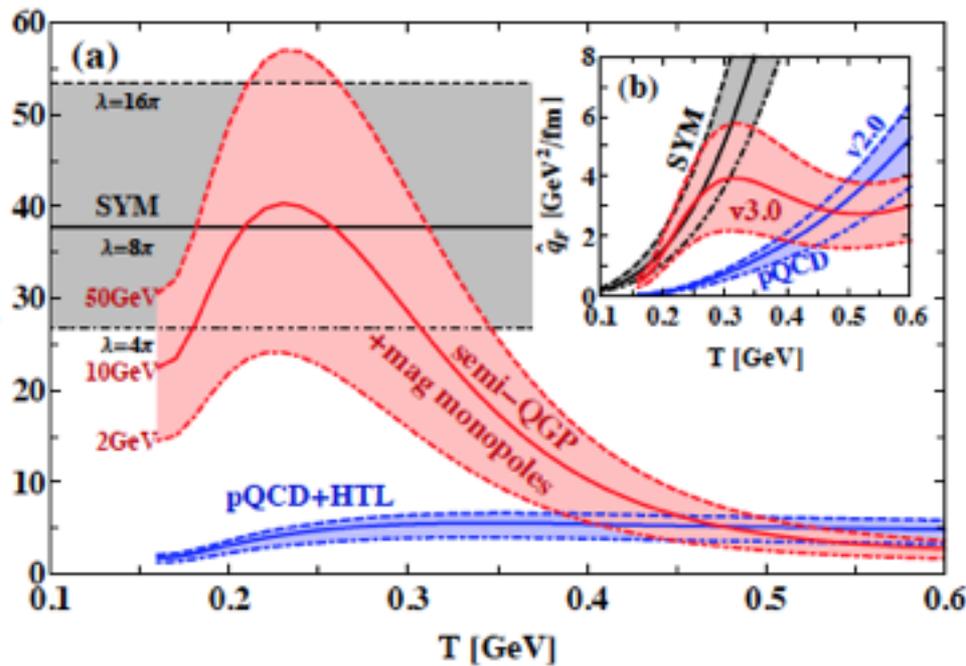


The SEVEN set of single hadron observables

$$\begin{aligned}
 & [ (RHIC+LHC) * (RAA+V2) * (pion) ] \\
 & + [ (LHC) * (RAA+V2) * (D) ] \\
 & + [ (LHC) * (RAA) * (B) ],
 \end{aligned}$$

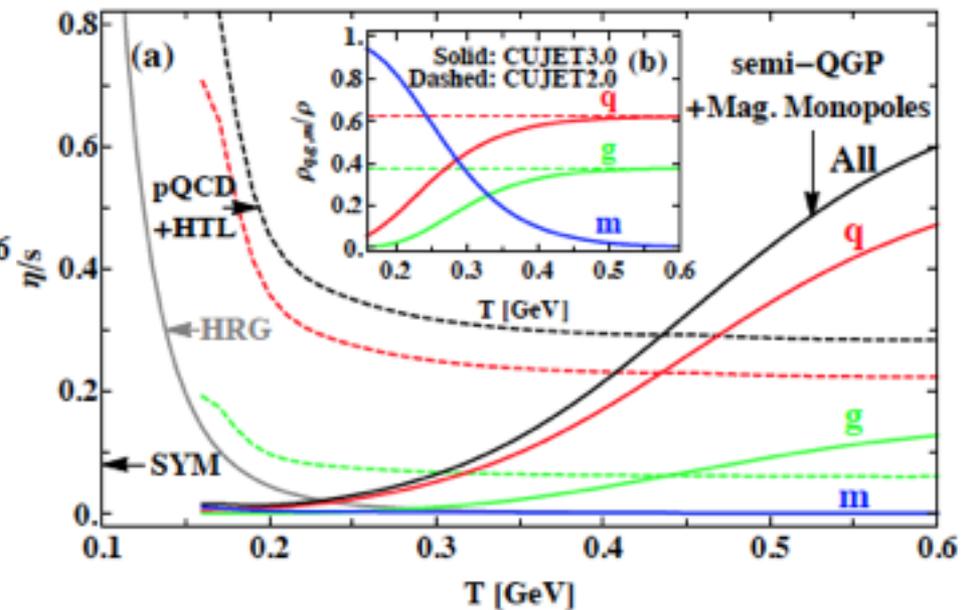
are nicely explained by CUJET3.0 framework  
(with essentially ONE model parameter)  
that implements the nonperturbative near- $T_c$  physics!

# Near- $T_c$ Matter Properties are Special!



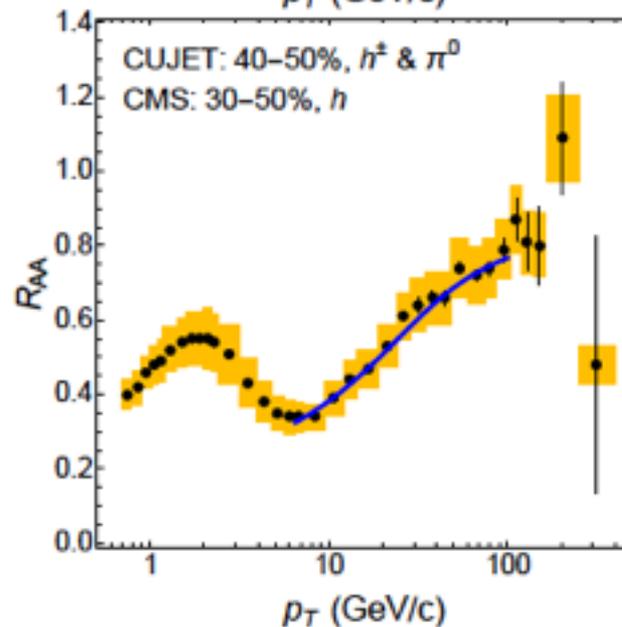
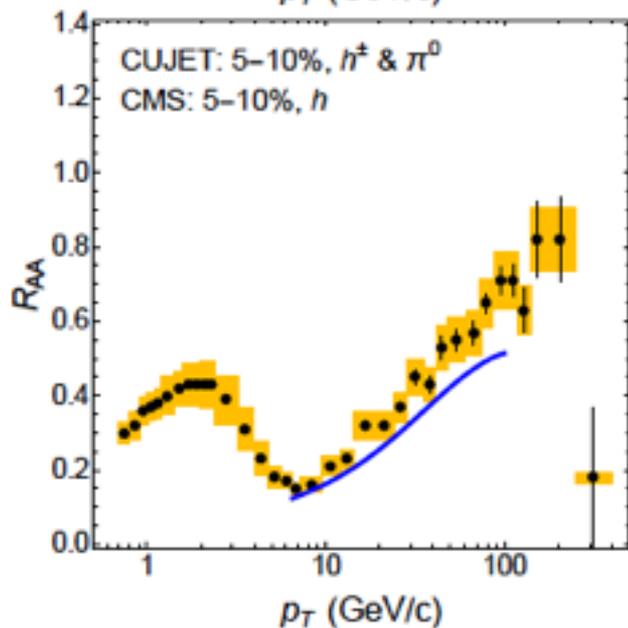
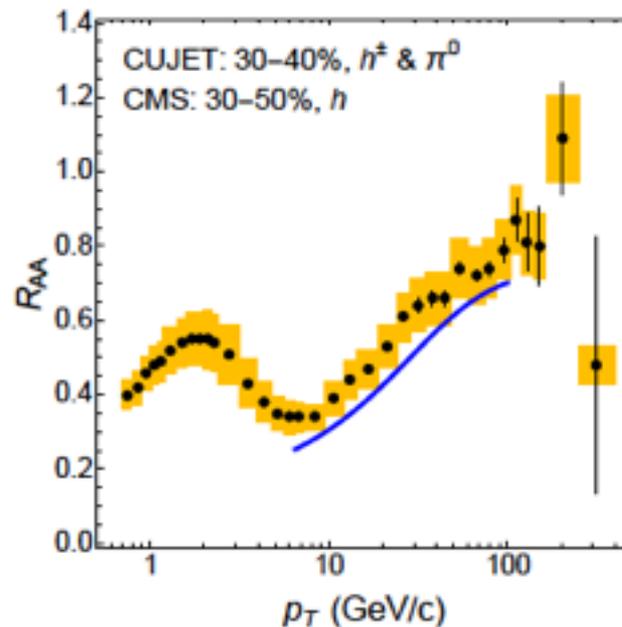
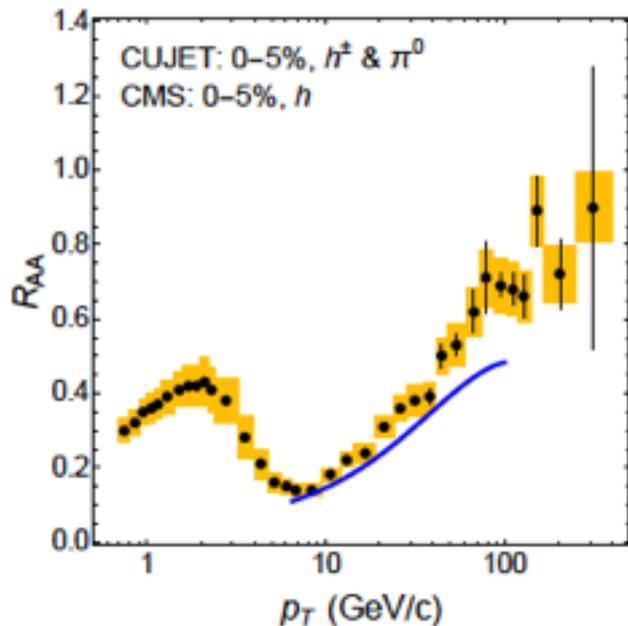
Jet transport coefficient  $\hat{q}/T^3$  computed from CUJET3.0 shows a prominent peak near  $T_c$ !

Shear viscosity, *etals*, computed from CUJET3.0 shows a clear minimum near  $T_c$  and rapid increase to high  $T$ !

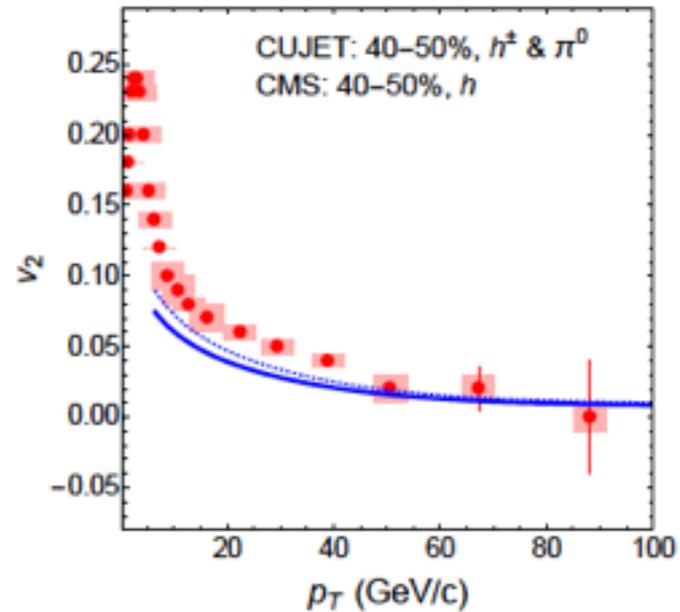
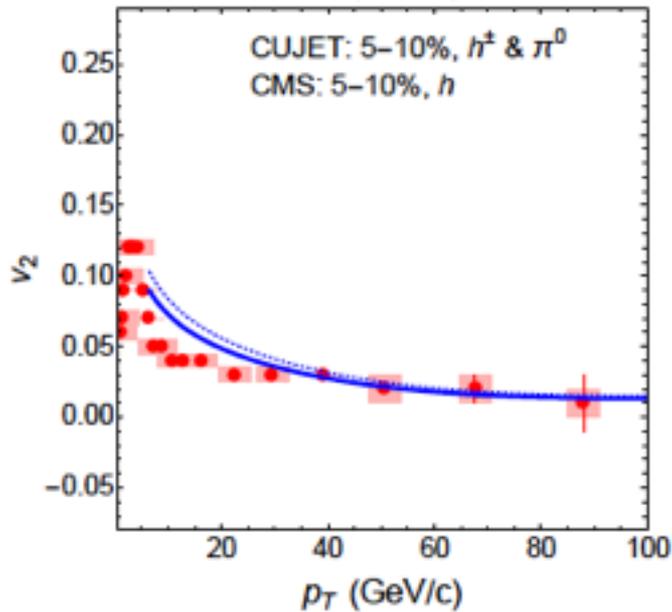
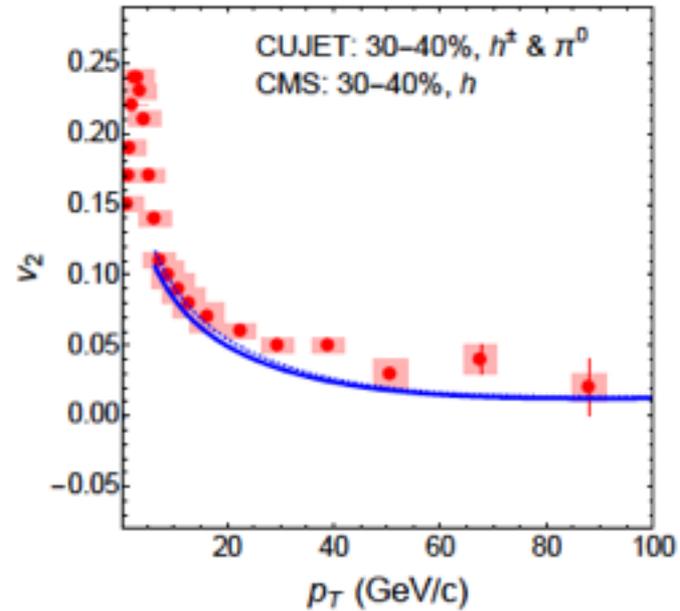
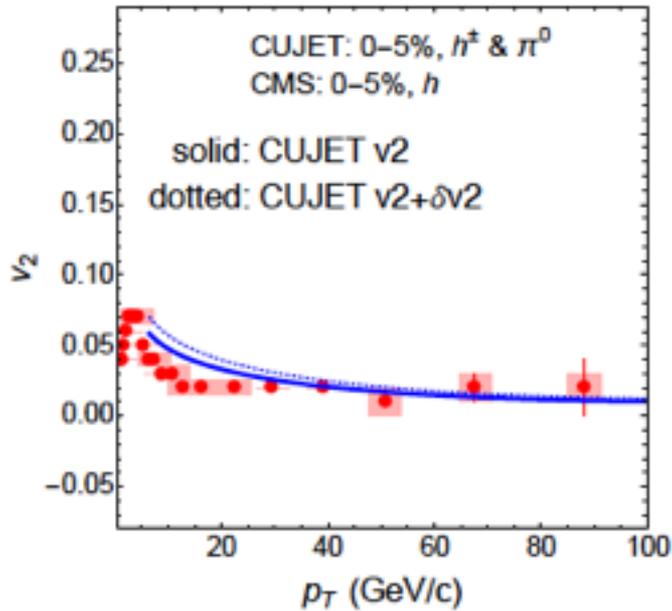


**CONSISTENCY of Perfect Fluidity & Jet Quenching in the semi-quark-gluon monopole plasma (sQGMP)!**

# Test CUJET3 @ LHC 5TeV

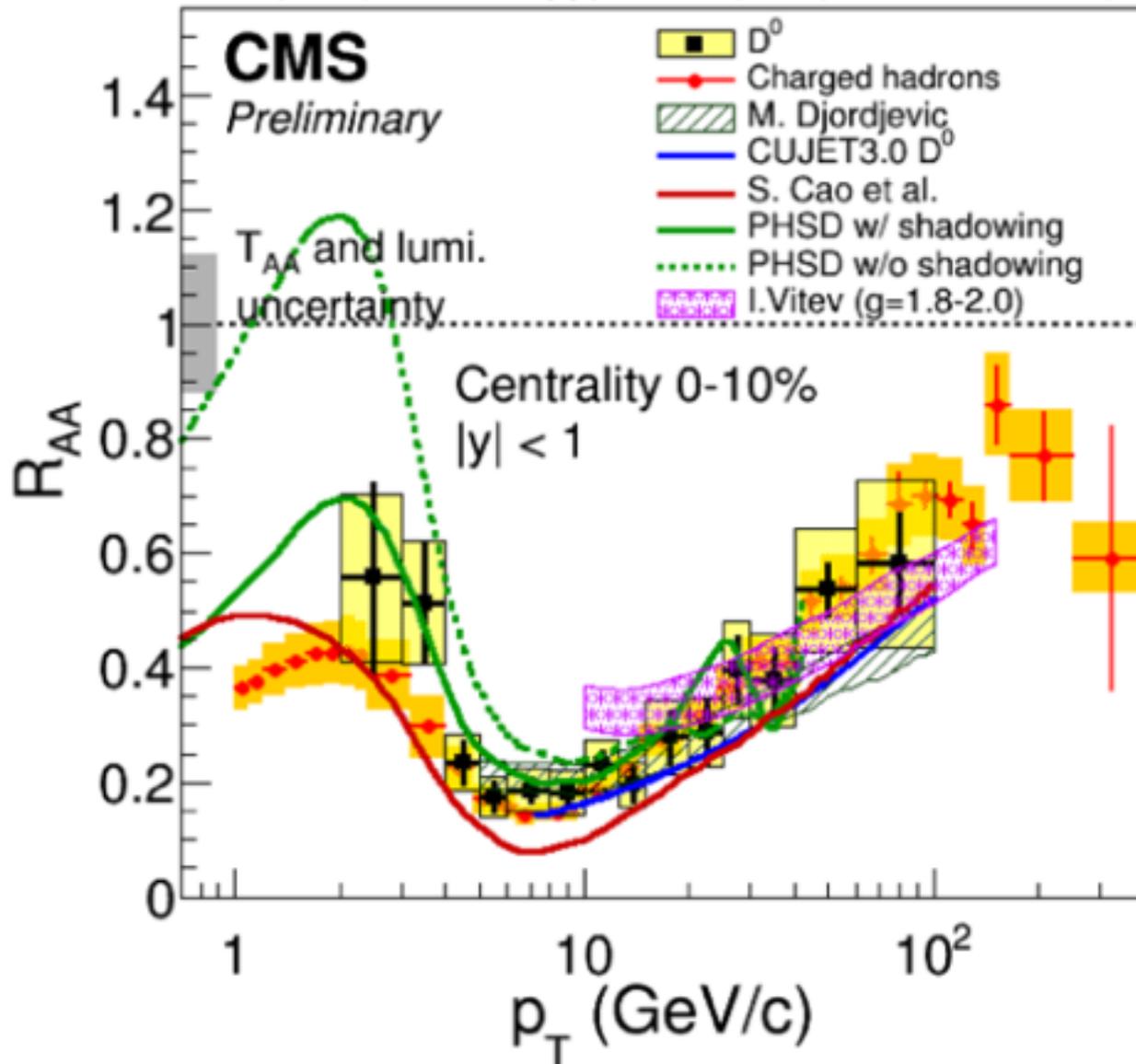


# Test CUJET3 @ LHC 5TeV

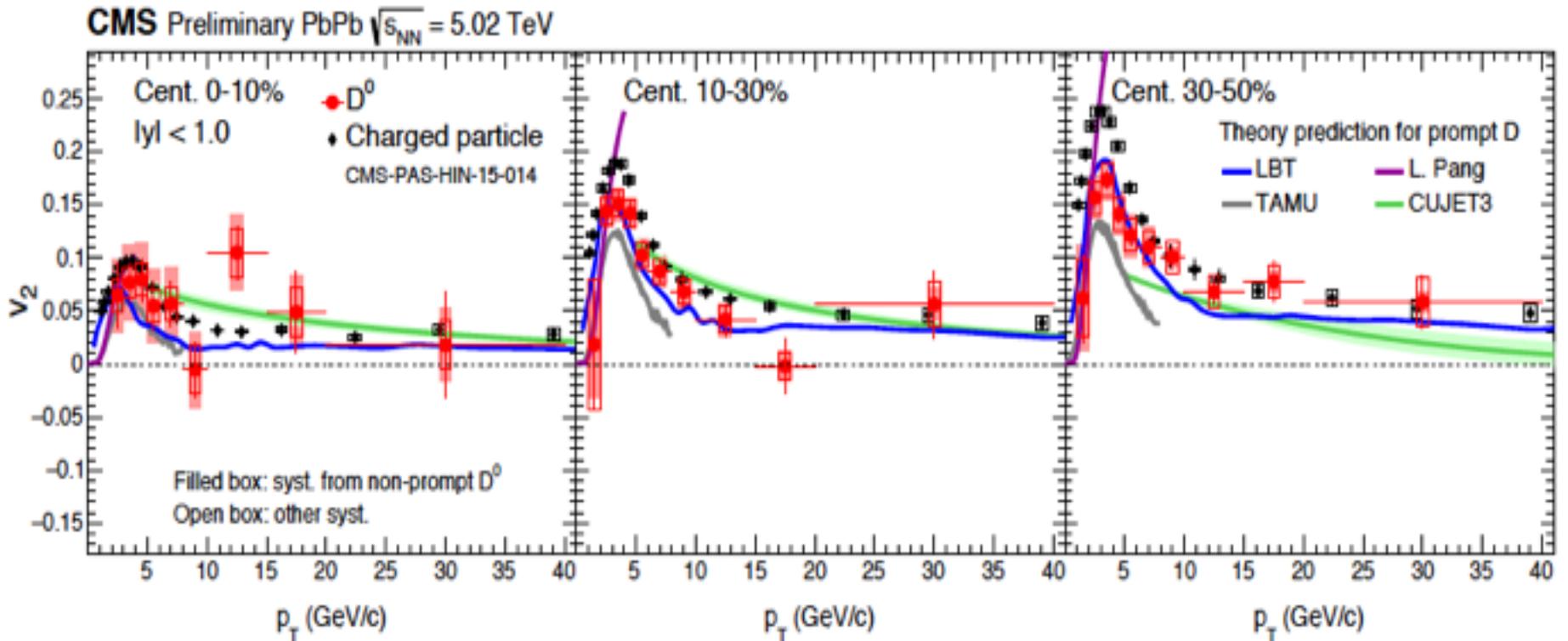


# Test CUJET3 @ LHC 5TeV

25.8 pb<sup>-1</sup> (5.02 TeV pp) + 404 μb<sup>-1</sup> (5.02 TeV PbPb)



# Test CUJET3 @ LHC 5TeV



# HF Test from CUJET3 for RRTF

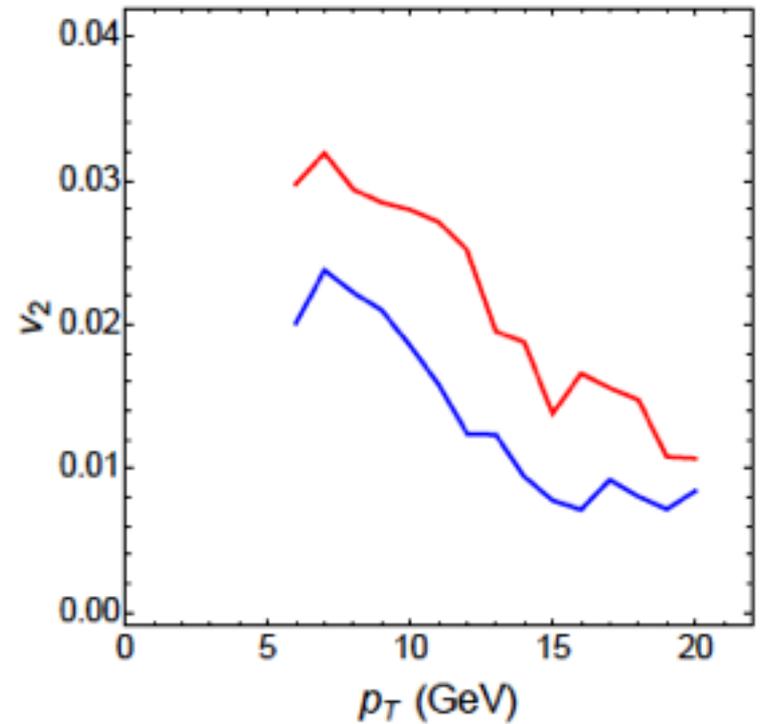
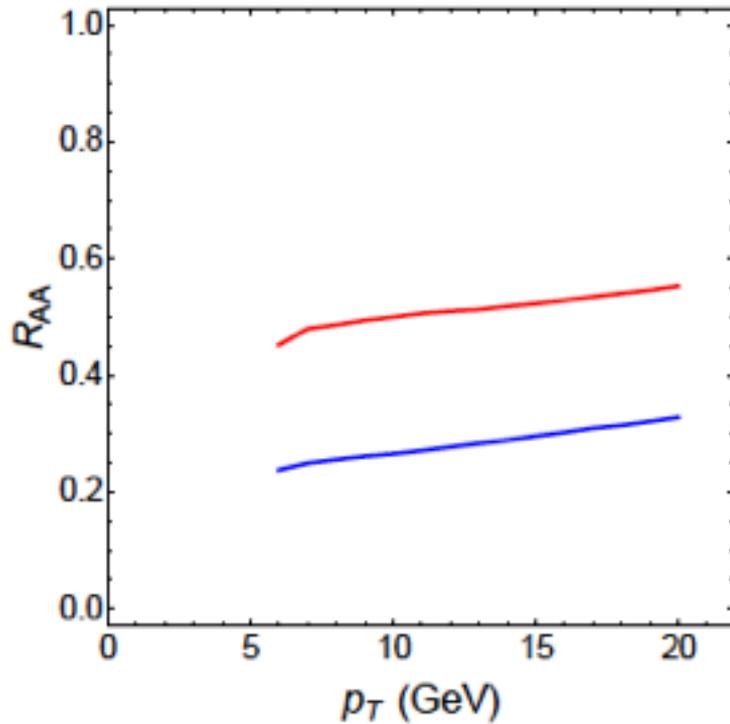
- Only collisional/elastic energy loss included.  
( no radiational/inelastic e.l. )
- $dE/dx = -A(p^2) p$  , [ M.Mustafa et al PRC57(1998) ]  
 $A(p^2)$  obtained from file “FPcoeff-pQCD-K5.dat”
- Both input charm quark pp spectrum & fragmentation function are using “RRTF standard”.
- 2+1D VISHNU hydro profile employed.
- For comparison, results using CUJET collisional energy loss kernel also attached.

$$\frac{dE}{dx} = -C_R \pi \alpha_s^2 T^2 \left( 1 + \frac{2}{6} \right) \left( \frac{1}{v} + \frac{v^2 - 1}{2v^2} \log \frac{1+v}{1-v} \right) \log \left( \frac{k_{max}}{\mu} \right)$$

# HF Test from CUJET3 for RRTF

$R_{AA}$  &  $v_2$  for charm quark

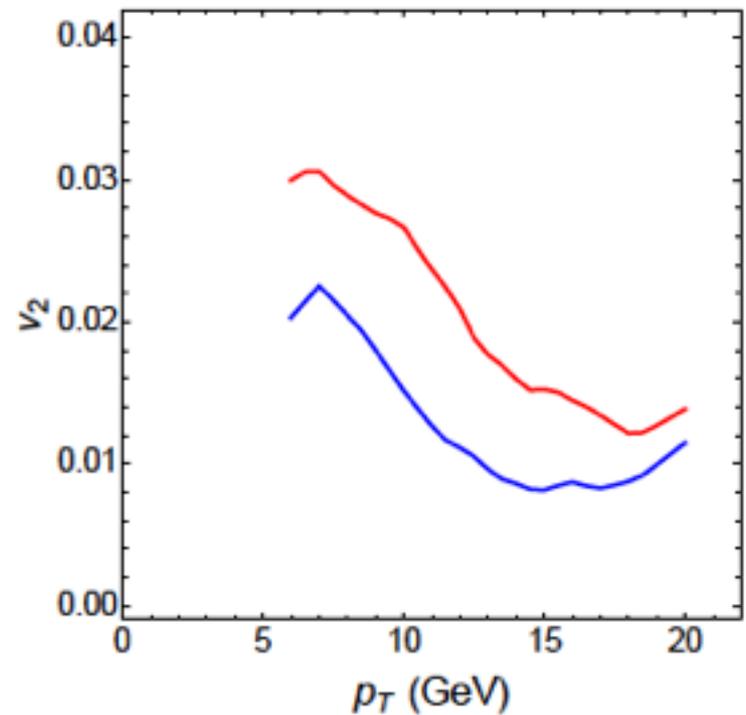
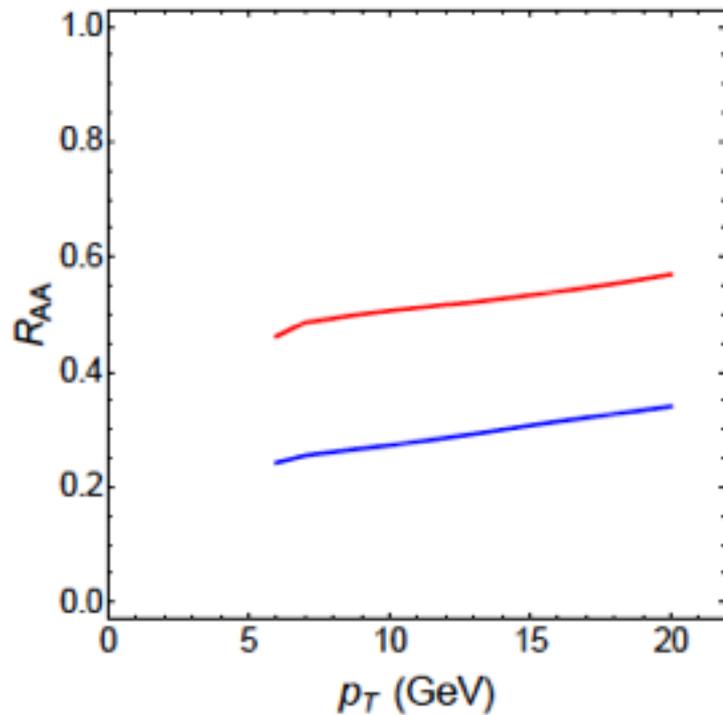
- Blue: 0-10%
- Red: 30-50%



# HF Test from CUJET3 for RRTF

$R_{AA}$  &  $v_2$  for D0

- Blue: 0-10%
- Red: 30-50%

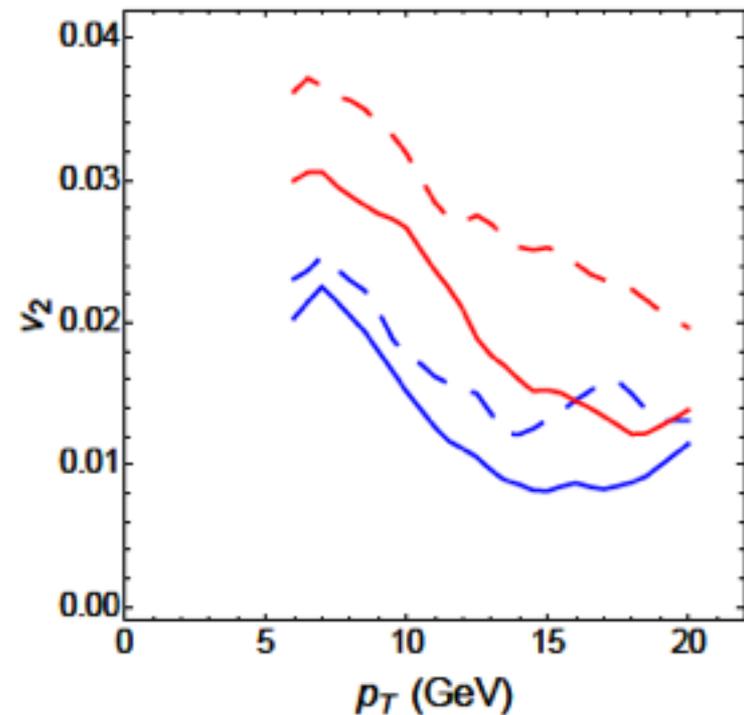
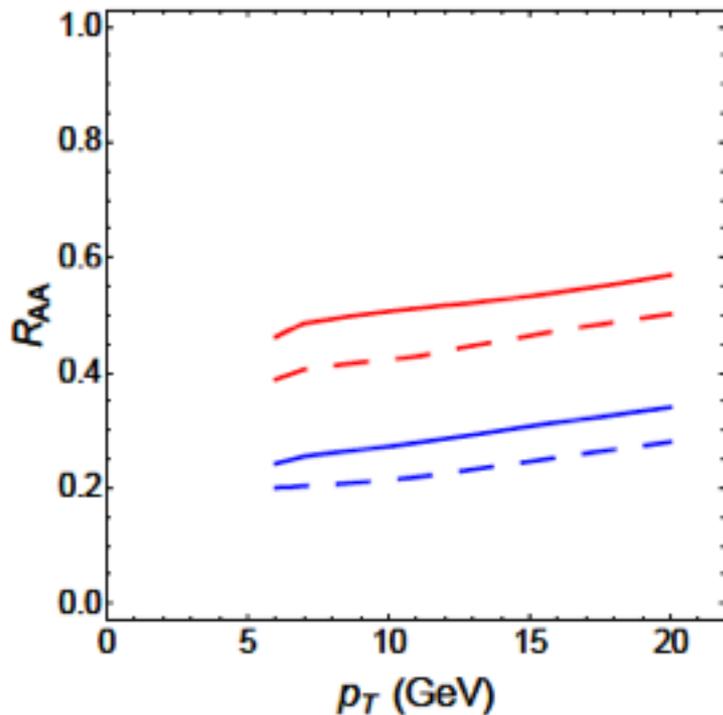


# HF Test from CUJET3 for RRTF

$R_{AA}$  &  $v_2$  for D0 — CUJET c.e.l. kernel

***for comparison***

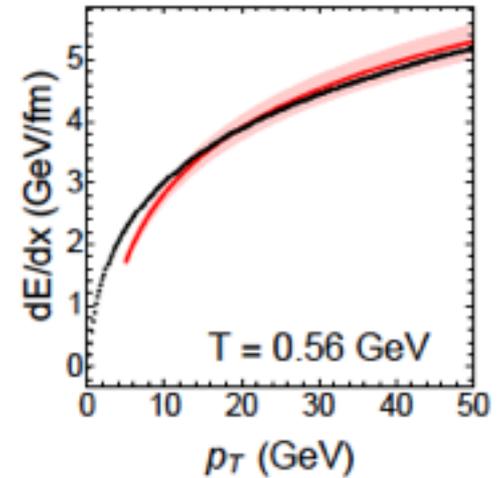
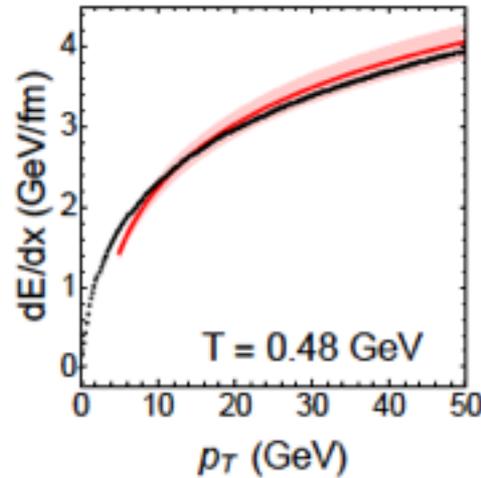
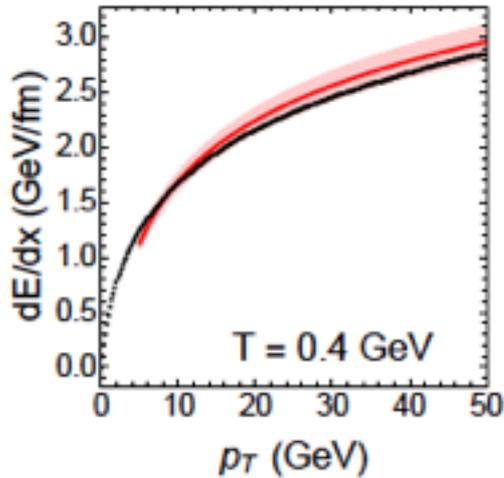
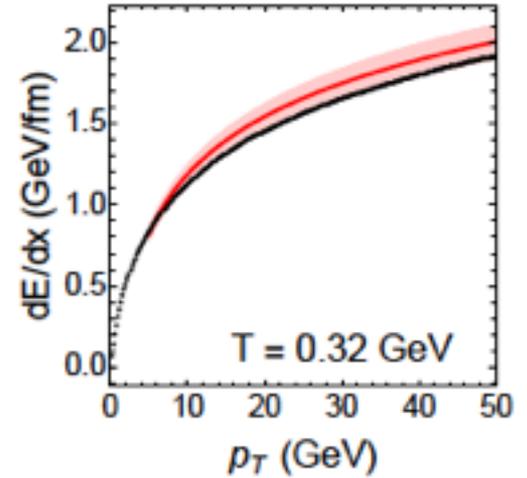
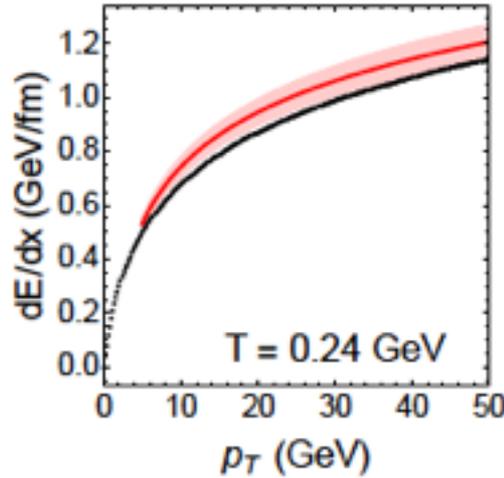
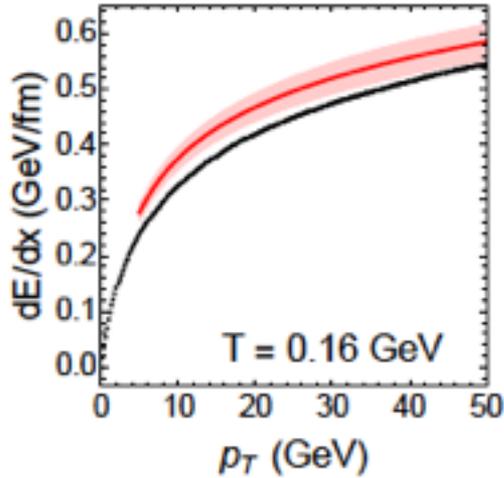
- Blue: 0-10%
- Red: 30-50%
- Solid: RRTF standard
- Dashing: CUJET with  $\alpha_s = 0.4$ ,  $N_f = 3$ ,  $m_c = 1.5$  GeV



# HF Test from CUJET3 for RRTF

Black dots: from "FPcoeff-pQCD-K5.dat"  
Red curves: CUJET standard kernel

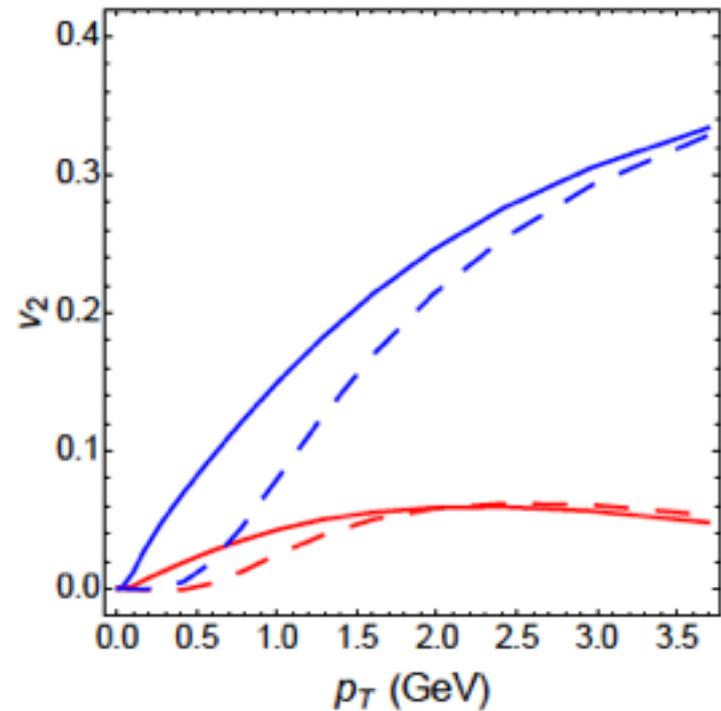
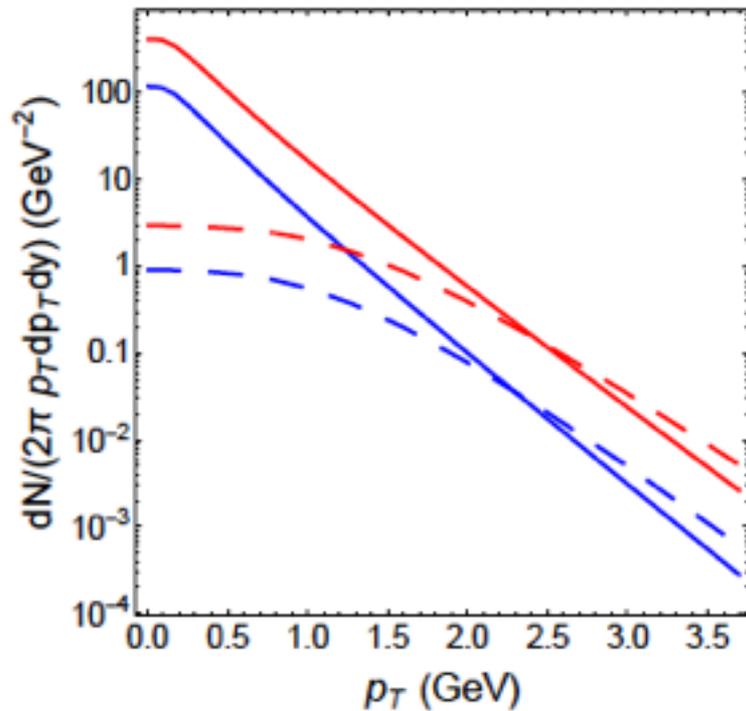
band:  $N_f = 2 \sim 3$  from bottom to top  
central curve:  $N_f = 2.5$



# HF Test from CUJET3 for RRTF

about hydro — Spectrum &  $v_2$  **BEFORE** hadron cascade  
(at the end of QGP evolution)

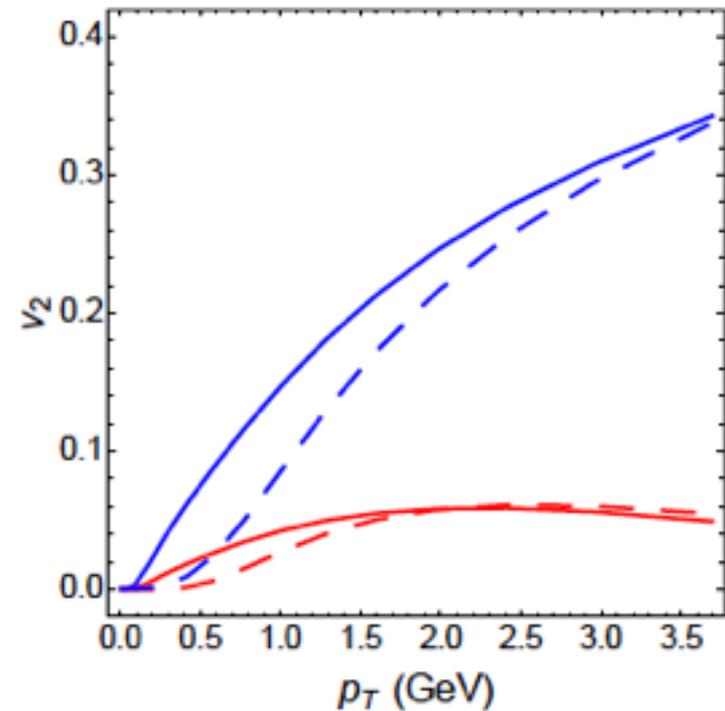
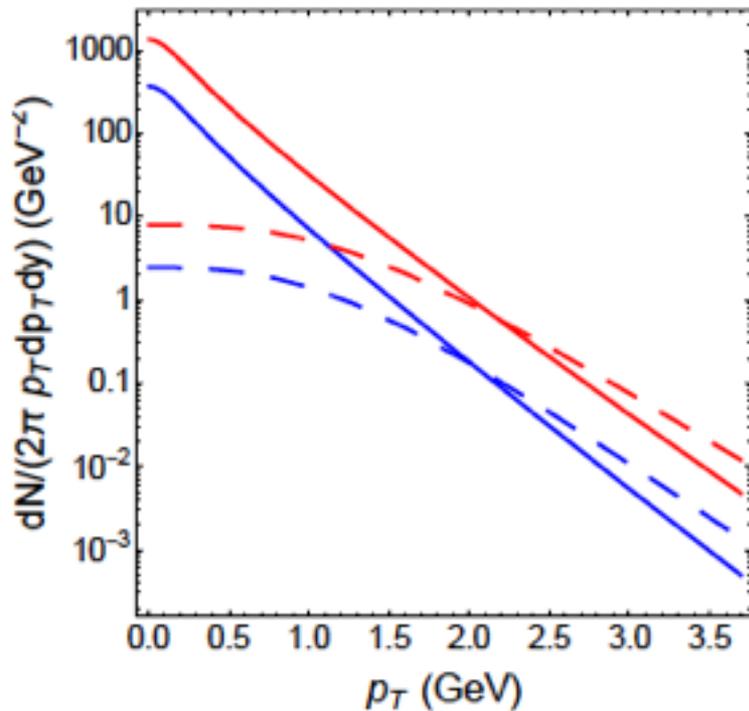
- Blue: 0-10%
- Red: 30-50%
- Solid: pions
- Dashed: protons



# HF Test from CUJET3 for RRTF

about hydro — Spectrum &  $v_2$  **AFTER** hadron cascade

- Blue: 0-10%
- Red: 30-50%
- Solid: pions
- Dashed: protons



# Summary & Discussions

